Theory overview of CLFV and EDM

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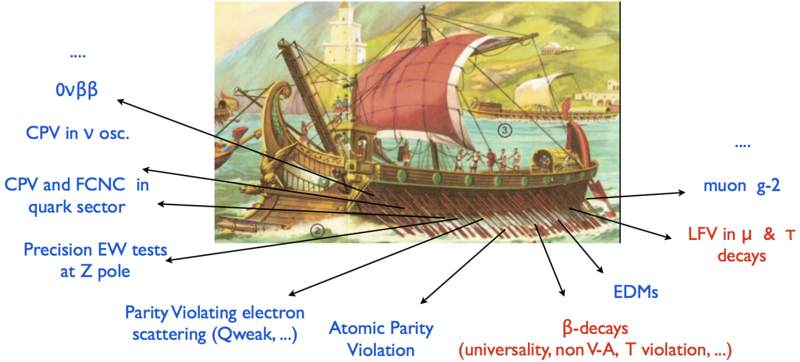
Plan of the talk

1. **Current status of LFV**
2. **LFV from heavy NP**
3. **EDMs, g-2 and LFV interrelationship**
4. **LFV @ FCC-ee & Muon Collider**
5. **LFV from light NP**
6. **Conclusions and future prospects**

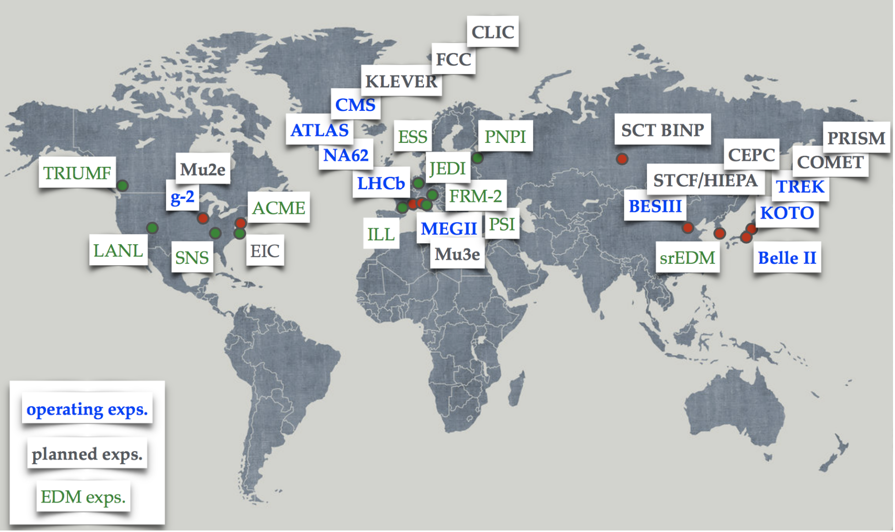
Where to look for New Physics at low-energy?

- **Processes very suppressed or even forbidden in the SM**
- **Processes predicted with high precision in the SM**



High-intensity frontier: A collective effort to determine the NP dynamics

Experimental status



Experimental bounds

Process	Present	Experiment	Future	Experiment
$\mu \to e\gamma$	1.5×10^{-13}	MEG	$\approx 6 \times 10^{-14}$	MEG II
$\mu \to 3e$	1.0×10^{-12}	SINDRUM	$\approx 10^{-16}$	Mu3e
μ^- Au $\to e^-$ Au	7.0×10^{-13}	SINDRUM II	?	
μ^- Ti $\to e^-$ Ti	4.3×10^{-12}	SINDRUM II	?	
μ^- Al $\to e^-$ Al	–		$\approx 10^{-16}$	COMET, MU2e
$\tau \to e\gamma$	3.3×10^{-8}	Belle & BaBar	$\sim 10^{-9}$	Belle II
$\tau \to \mu\gamma$	4.4×10^{-8}	Belle & BaBar	$\sim 10^{-9}$	Belle II
$\tau \to 3e$	2.7×10^{-8}	Belle & BaBar	$\sim 10^{-10}$	Belle II
$\tau \to 3\mu$	2.1×10^{-8}	Belle & BaBar	$\sim 10^{-10}$	Belle II
$d_e(\mathrm{e\ cm})$	1.1×10^{-29}	ACME	$\sim 3 \times 10^{-31}$	ACME III
$d_\mu(\mathrm{e\ cm})$	1.8×10^{-19}	Muon (g-2)	$\sim 10^{-22}$	PSI

Table: Present and future experimental sensitivities for relevant low-energy observables.

- So far, only upper bounds. Still excellent prospects for exp. improvements.
- We can expect a NP signal in all above observables below the current bounds.

Charged LFV in the SM

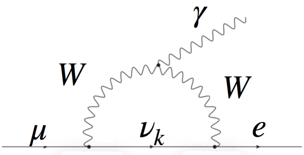
- GIM mechanism very effective in LFV transitions
- amplitude proportional to $A(\mu \to e\gamma) \propto m_\nu^2$

Very small !!!

$$\mathrm{BR}(\mu \to e\gamma) \simeq \frac{\Gamma(\mu \to e\gamma)}{\Gamma(\mu \to e\nu\bar{\nu})} = \frac{3\alpha}{32\pi} \left| \sum_{k=1,3} \frac{U_{\mu k} U^*_{ek} m^2_{\nu_k}}{M_W^2} \right|^2 .$$

$$\mathrm{BR}(\mu \to e\gamma) = 10^{-55} \div 10^{-54}$$

- similar suppressions for $\mu \to 3e, \tau \to 3\mu, \mu \to e, \ldots$

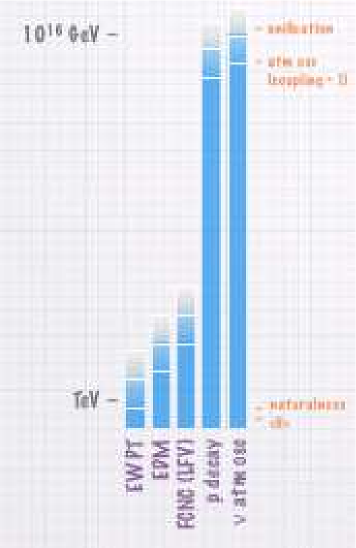


Why flavor violation is visible in neutrino oscillation while it's not in charged LFV?
The uncertainty principle sets the oscillation time for $\mu \to e\gamma$ to be $t \sim h/M_W$!

Message: Any evidence for LFV would be an unambiguous signal of NP!

Why do we need New Physics (NP)?

- **Gravity** $\Longrightarrow$ $\Lambda_{\text{Planck}} \sim 10^{18-19}$ GeV
- **Neutrino masses** $\Longrightarrow$ $\Lambda_{\text{see-saw}} \lesssim 10^{15}$ GeV
- **BAU**: evidence of CPV beyond SM
  - Electroweak Baryogenesis $\Longrightarrow$ $\Lambda_{\text{NP}} \lesssim$ TeV
  - Leptogenesis $\Longrightarrow$ $\Lambda_{\text{see-saw}} \lesssim 10^{15}$ GeV
- **Dark Matter (WIMP)** $\Longrightarrow$ $\Lambda_{\text{NP}} \lesssim$ TeV
- **Hierarchy problem**: $\Longrightarrow$ $\Lambda_{\text{NP}} \lesssim$ TeV

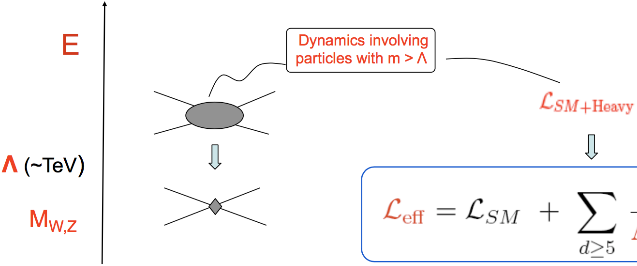


SM = effective theory at the EW scale

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_{d\geq 5} \frac{c_{ij}^{(d)}}{\Lambda_{NP}^{d-4}}\, O_{ij}^{(d)}$$

- $\mathcal{L}_{\text{eff}}^{d=5} = \frac{y_\nu^{ij}}{\Lambda_{\text{see-saw}}} L_i L_j \phi\phi$,
- $\mathcal{L}_{\text{eff}}^{d=6}$ generates FCNC operators

- **Dynamics below the scale Λ [$\sim$ mass of new particles] is described by L_{eff}**

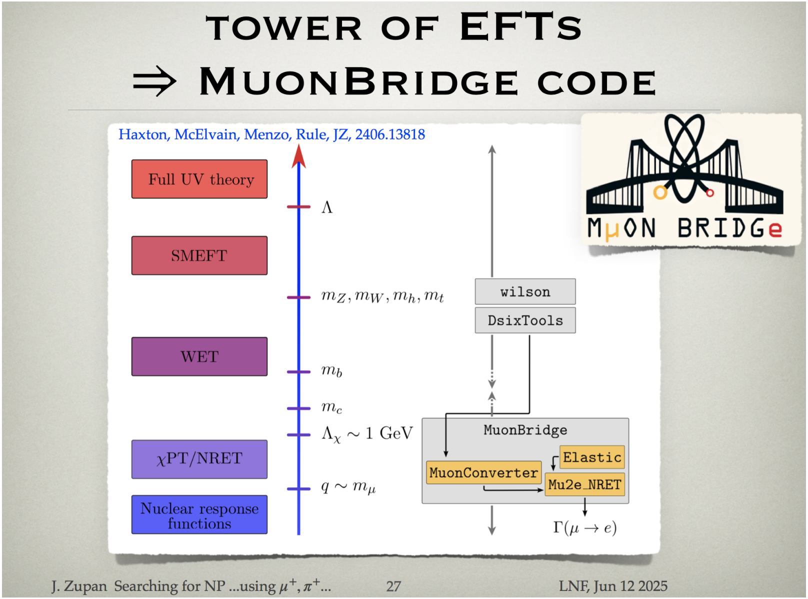


$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{SM} + \sum_{d \geq 5} \frac{c_n^{(d)}}{\Lambda^{d-4}} \hat{O}_n^{(d)}[\phi_{SM}]$$

  - L_{eff} is built out of relevant low-energy degrees of freedom (SM fields)
  - L_{eff} respects the SM gauge symmetries $G_{\text{SM}} = SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$
  - L_{eff} is organized in inverse powers of Λ (amplitudes suppressed by powers of E/Λ)

- **Experiments at the precision frontier probe energy scale Λ and symmetries of the new interactions (coeff. & structure of $\hat{O}_n^{(d)}$)**

MuonBridge: a $\mu \to e$ conversion in Nuclei code

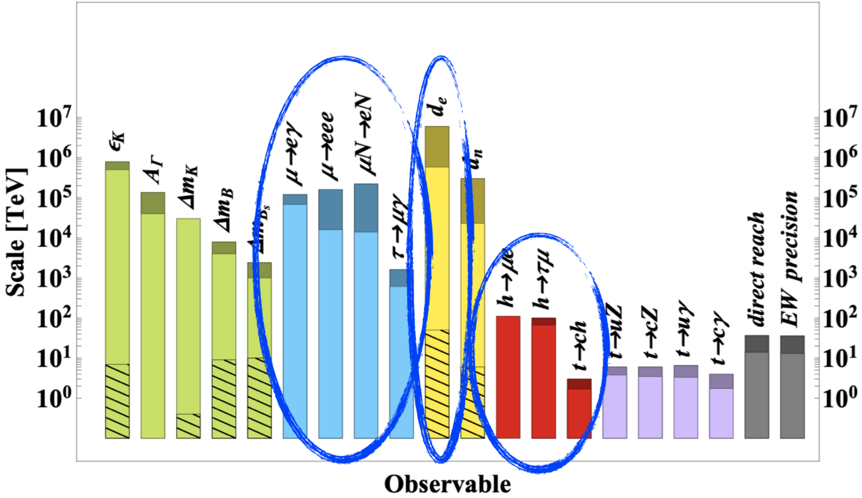


EFT approach to LFV

4-leptons operators		Dipole operators	
$Q_{\ell\ell}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{L}_L\gamma^\mu L_L)$	Q_{eW}	$(\bar{L}_L\sigma^{\mu\nu}e_R)\tau_I\Phi W^I_{\mu\nu}$
Q_{ee}	$(\bar{e}_R\gamma_\mu e_R)(\bar{e}_R\gamma^\mu e_R)$	Q_{eB}	$(\bar{L}_L\sigma^{\mu\nu}e_R)\Phi B_{\mu\nu}$
$Q_{\ell e}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{e}_R\gamma^\mu e_R)$		
2-lepton 2-quark operators			
$Q^{(1)}_{\ell q}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{Q}_L\gamma^\mu Q_L)$	$Q_{\ell u}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{u}_R\gamma^\mu u_R)$
$Q^{(3)}_{\ell q}$	$(\bar{L}_L\gamma_\mu\tau_I L_L)(\bar{Q}_L\gamma^\mu\tau_I Q_L)$	Q_{eu}	$(\bar{e}_R\gamma_\mu e_R)(\bar{u}_R\gamma^\mu u_R)$
Q_{eq}	$(\bar{e}_R\gamma^\mu e_R)(\bar{Q}_L\gamma_\mu Q_L)$	$Q_{\ell edq}$	$(\bar{L}^a_L e_R)(\bar{d}_R Q^a_L)$
$Q_{\ell d}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{d}_R\gamma^\mu d_R)$	$Q^{(1)}_{\ell equ}$	$(\bar{L}^a_L e_R)\epsilon_{ab}(\bar{Q}^b_L u_R)$
Q_{ed}	$(\bar{e}_R\gamma_\mu e_R)(\bar{d}_R\gamma^\mu d_R)$	$Q^{(3)}_{\ell equ}$	$(\bar{L}^a_i\sigma_{\mu\nu}e_R)\epsilon_{ab}(\bar{Q}^b_L\sigma^{\mu\nu}u_R)$
Lepton-Higgs operators			
$Q^{(1)}_{\Phi\ell}$	$(\Phi^\dagger i \overleftrightarrow{D}_\mu \Phi)(\bar{L}_L\gamma^\mu L_L)$	$Q^{(3)}_{\Phi\ell}$	$(\Phi^\dagger i \overleftrightarrow{D}^I_\mu \Phi)(\bar{L}_L\tau_I\gamma^\mu L_L)$
$Q_{\Phi e}$	$(\Phi^\dagger i \overleftrightarrow{D}_\mu \Phi)(\bar{e}_R\gamma^\mu e_R)$	$Q_{e\Phi 3}$	$(\bar{L}_L e_R\Phi)(\Phi^\dagger\Phi)$

$\mu \to e\gamma$ $\mu \to 3e$ $\mu \to e$

Bounds on the NP scale



[Physics Briefing Book, 1910.11775]

Correlating LFV signals

- **LFV** operators **@ dim-6**

$$\mathcal{L}_{\rm eff} = \mathcal{L}_{\rm SM} + \frac{1}{\Lambda^2_{\rm LFV}}\,\mathcal{O}^{\rm dim-6} + \ldots\,.$$

$$\mathcal{O}^{\rm dim-6} \ni \;\; \bar\mu_R\,\sigma^{\mu\nu}\,H\,e_L\,F_{\mu\nu}\,, \;\; (\bar\mu_L\gamma^\mu e_L)\left(\bar f_L\gamma^\mu f_L\right)\,, \;\; (\bar\mu_R e_L)\left(\bar f_R f_L\right)\,, \; f = e, u, d$$

- $\ell \to \ell' \gamma$ probe ONLY the dipole-operator (at tree level)
- $\ell_i \to \ell_j \bar\ell_k \ell_k$ and $\mu \to e$ in Nuclei probe dipole and 4-fermion operators
- When the dipole-operator is dominant:

$$\mathrm{BR}(\ell_i \to \ell_j \ell_k \bar\ell_k) \approx \alpha \times \mathrm{BR}(\ell_i \to \ell_j \gamma)$$
$$\mathrm{CR}(\mu \to e \text{ in N}) \approx \alpha \times \mathrm{BR}(\mu \to e\gamma)$$

$$\frac{\mathrm{BR}(\mu \to \mathbf{3e})}{\mathbf{3 \times 10^{-15}}} \approx \frac{\mathrm{BR}(\mu \to \mathbf{e}\gamma)}{\mathbf{5 \times 10^{-13}}} \approx \frac{\mathrm{CR}(\mu \to \mathbf{e} \text{ in N})}{\mathbf{3 \times 10^{-15}}}$$

- **Ratios like** $Br(\mu \to e\gamma)/Br(\tau \to \mu\gamma)$ **probe the NP flavor structure**
- **Ratios like** $Br(\mu \to e\gamma)/Br(\mu \to eee)$ **probe the NP operator at work**

Polarising the muon to distinguish operators

Kuno, Okada hep-ph/9909265

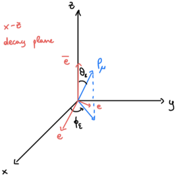
$$\frac{dB(\mu \to e\gamma)}{d(\cos\theta_e)} = 192\pi^2 \left(\frac{v}{\Lambda}\right)^4 \left[\left|C_{D,R}\right|^2 \left(1 - P_\mu \cos\theta_e\right) + \left|C_{D,L}\right|^2 \left(1 + P_\mu \cos\theta_e\right) \right]$$

Muon polarization vector — **Angle between e momentum and $\vec{P}_\mu$**

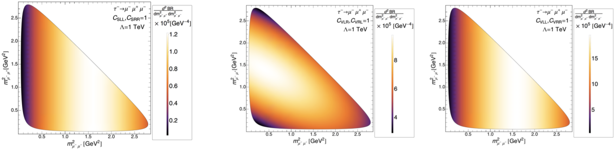
Petcov, Bolton 2204.03468

$$\frac{dB_{\mu\to 3e}}{dx_1 dx_2 d\Omega_e} = \frac{3}{2\pi}\left[C_1(x_1,x_2) + C_2(x_1,x_2)\, P_\mu \cos\theta_e + C_3(x_1,x_2)\, P_\mu \sin\theta_e \cos\phi_e + C_4(x_1,x_2)\, P_\mu \sin\theta_e \sin\phi_e\right]$$

Can distinguish $C_{V,LX}, C_{V,LX}, C_{S,R}$ from $C_{V,RX}, C_{V,RX}, C_{S,L}$ but not scalars from vectors
CP asymmetries are also measureable (phase between dipoles and vectors)

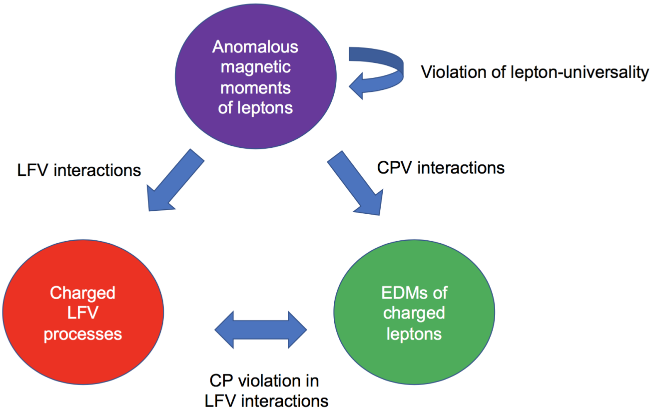


- Dalitz plots for the three body also possible to distinguish operators (vector vs scalar)



Scalars **Vectors**

Probing NP through leptonic dipole moments



On leptonic dipoles: $\ell \to \ell'\gamma$

- **NP effects are encoded in the effective Lagrangian**

$$\mathcal{L} = e\frac{m_\ell}{2}\left(\bar{\ell}_R\sigma_{\mu\nu}A_{\ell\ell'}\ell'_L + \bar{\ell}'_L\sigma_{\mu\nu}A^\star_{\ell\ell'}\ell_R\right)F^{\mu\nu} \qquad \ell,\ell' = e,\mu,\tau\,,$$

  - **Branching ratios of $\ell \to \ell'\gamma$**

$$\frac{\mathrm{BR}(\ell\to\ell'\gamma)}{\mathrm{BR}(\ell\to\ell'\nu_\ell\bar{\nu}_{\ell'})} = \frac{48\pi^3\alpha}{G_F^2}\left(|A_{\ell\ell'}|^2 + |A_{\ell'\ell}|^2\right).$$

  - **Δa_ℓ and leptonic EDMs**

$$\Delta a_\ell = 2m_\ell^2\,\mathrm{Re}(A_{\ell\ell}), \qquad \frac{d_\ell}{e} = m_\ell\,\mathrm{Im}(A_{\ell\ell})\,.$$

  - **"Naive scaling":** a broad class of NP theories contributes to Δa_ℓ and d_ℓ as

$$\frac{\Delta a_\ell}{\Delta a_{\ell'}} = \frac{m_\ell^2}{m_{\ell'}^2}, \qquad \frac{d_\ell}{d_{\ell'}} = \frac{m_\ell}{m_{\ell'}}\,.$$

- $\mathrm{BR}(\ell_i \to \ell_j\gamma)$ **vs.** $(g-2)_\mu$

$$\mathrm{BR}(\mu \to e\gamma) \approx 3 \times 10^{-13} \left(\frac{\Delta a_\mu}{3 \times 10^{-9}}\right)^2 \left(\frac{\theta_{e\mu}}{10^{-5}}\right)^2$$

$$\mathrm{BR}(\tau \to \mu\gamma) \approx 4 \times 10^{-8} \left(\frac{\Delta a_\mu}{3 \times 10^{-9}}\right)^2 \left(\frac{\theta_{\mu\tau}}{10^{-2}}\right)^2$$

- **EDMs vs.** $(g-2)_\mu$

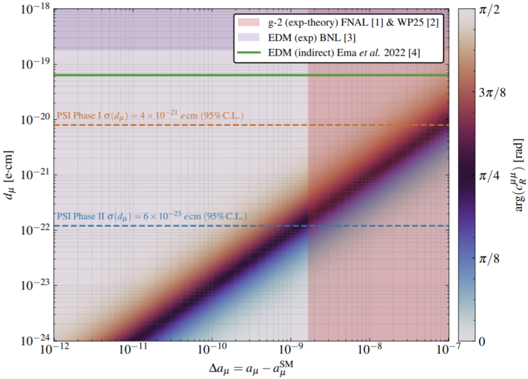
$$d_e \simeq \left(\frac{\Delta a_\mu}{3 \times 10^{-9}}\right) 10^{-29} \left(\frac{\phi_e^{CPV}}{10^{-5}}\right) \; e\,\mathrm{cm}\,,$$

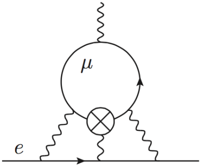
$$d_\mu \simeq \left(\frac{\Delta a_\mu}{3 \times 10^{-9}}\right) 2 \times 10^{-22} \, \phi_\mu^{CPV} \; e\,\mathrm{cm}\,,$$

- **Main messages:**
  - $\Delta a_\mu = 38(63) \times 10^{-11}$ **requires a nearly flavor and CP conserving NP**
  - **Large effects in the muon EDM** $d_\mu \sim 10^{-22}\, e\,\mathrm{cm}$ **are still allowed!**

[Giudice, P.P., & Passera, '12]

Experimental status of the muon EDM





$$d_\mu \leq 10^{-21}\, e\,\text{cm} \left(\frac{d_e}{10^{-31}\, e\,\text{cm}}\right)$$

[Crivellin, Hoferichter & Schmidt-Wellenburg, '18]

$$d_\mu \simeq \left(\frac{\Delta a_\mu}{3\times10^{-9}}\right) 2\times10^{-22}\, \phi_\mu^{CPV}\ e\,\text{cm}\,,$$

[Giudice, PP & Passera, '12]

Low-energy vs high-energy LFV

Purely leptonic contact interations:
- $e^+e^- \to \mu\tau$,
- $\mu^+\mu^- \to \mu e$,
- $\mu^- \to eee$, … .

4-leptons operators	
$Q_{\ell\ell}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{L}_L\gamma^\mu L_L)$
Q_{ee}	$(\bar{e}_R\gamma_\mu e_R)(\bar{e}_R\gamma^\mu e_R)$
$Q_{\ell e}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{e}_R\gamma^\mu e_R)$

Dipole operators	
Q_{eW}	$(\bar{L}_L\sigma^{\mu\nu}e_R)\tau_I\Phi W^I_{\mu\nu}$
Q_{eB}	$(\bar{L}_L\sigma^{\mu\nu}e_R)\Phi B_{\mu\nu}$

Loop generated dipoles:
- $\mu^- \to e\gamma$,
- $\tau^- \to e\gamma$,
- $\tau^- \to \mu\gamma$.

2-lepton 2-quark operators			
$Q^{(1)}_{\ell q}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{Q}_L\gamma^\mu Q_L)$	$Q_{\ell u}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{u}_R\gamma^\mu u_R)$
$Q^{(3)}_{\ell q}$	$(\bar{L}_L\gamma_\mu\tau_I L_L)(\bar{Q}_L\gamma^\mu\tau_I Q_L)$	Q_{eu}	$(\bar{e}_R\gamma_\mu e_R)(\bar{u}_R\gamma^\mu u_R)$
Q_{eq}	$(\bar{e}_R\gamma^\mu e_R)(\bar{Q}_L\gamma_\mu Q_L)$	$Q_{\ell edq}$	$(\bar{L}^a_L e_R)(\bar{d}_R Q^a_L)$
$Q_{\ell d}$	$(\bar{L}_L\gamma_\mu L_L)(\bar{d}_R\gamma^\mu d_R)$	$Q^{(1)}_{\ell equ}$	$(\bar{L}^a_L e_R)\epsilon_{ab}(\bar{Q}^b_L u_R)$
Q_{ed}	$(\bar{e}_R\gamma_\mu e_R)(\bar{d}_R\gamma^\mu d_R)$	$Q^{(3)}_{\ell equ}$	$(\bar{L}^a_i\sigma_{\mu\nu}e_R)\epsilon_{ab}(\bar{Q}^b_L\sigma^{\mu\nu}u_R)$

Semi-leptonic operators:
- $\mu N \to e N$,
- $pp \to l_i l_k$,
- $B \to \pi e\mu$,
- $B \to K\mu\tau$, … .

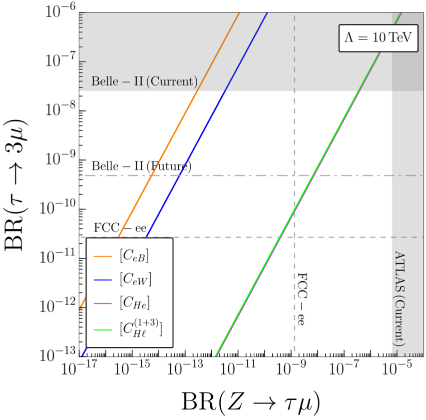
Lepton-Higgs Interactions:
- $h \to \tau\mu$,
- $Z \to \tau\mu$,
- $e^+e^- \to e\tau$,
- $\tau^- \to \mu ee$.

Lepton-Higgs operators			
$Q^{(1)}_{\Phi\ell}$	$(\Phi^\dagger i \overleftrightarrow{D}_\mu \Phi)(\bar{L}_L\gamma^\mu L_L)$	$Q^{(3)}_{\Phi\ell}$	$(\Phi^\dagger i \overleftrightarrow{D}^I_\mu \Phi)(\bar{L}_L\tau_I\gamma^\mu L_L)$
$Q_{\Phi e}$	$(\Phi^\dagger i \overleftrightarrow{D}_\mu \Phi)(\bar{e}_R\gamma^\mu e_R)$	$Q_{e\Phi 3}$	$(\bar{L}_L e_R\Phi)(\Phi^\dagger\Phi)$

Energy dependence of LFV signals

$$\Gamma(\tau \to 3e) \sim \frac{m_\tau^5}{\Lambda^4} \qquad \Gamma(Z \to \tau e) \sim \frac{m_Z^5}{\Lambda^4} \qquad \sigma(ee \to \tau e) \sim \frac{E_{CM}^2}{\Lambda^4}$$

Message: Interplay between high- and low-energy/high-statistics frontiers!



[Allwicher, Bartocci, Cornella, Leal, Paradisi, Scantamburlo, Sumensari, to appear.]

	Current	Future	FCC-ee
$\tau \to e\gamma$	3.3×10^{-8}	1.2×10^{-8}	–
$\tau \to \mu\gamma$	4.2×10^{-8}	5×10^{-9}	1.2×10^{-9}
$\tau \to eee$	2.7×10^{-8}	4×10^{-10}	–
$\tau \to \mu\mu\mu$	1.9×10^{-8}	3×10^{-10}	2.0×10^{-11}
$Z \to e\tau$	5.0×10^{-6}	–	10^{-9}
$Z \to \mu\tau$	6.5×10^{-6}	–	10^{-9}
$h \to e\tau$	2.0×10^{-3}	2.4×10^{-4}	1.6×10^{-4}
$h \to \mu\tau$	1.5×10^{-3}	2.4×10^{-4}	1.4×10^{-4}

[FCC-ee projections by Lusiani et al.]

Large statistics expected at FCC-ee:
- $\sim 10^{12}$ Z bosons
- $\sim 10^{11}$ τ pairs

Message: LFV tau decays are more powerful than Z decays at FCC-ee!

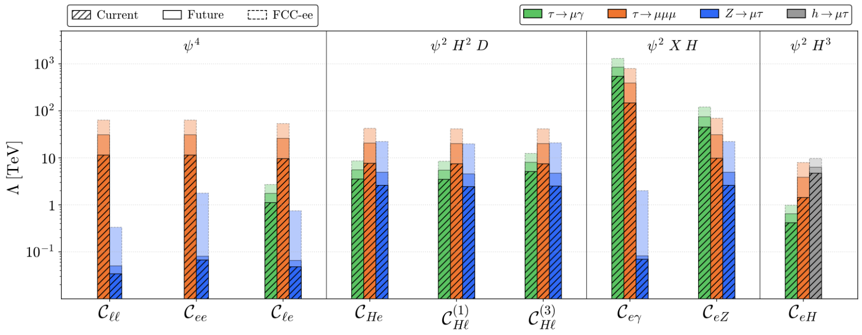


FIG. 2. *Reach in the new-physics scale Λ (for $\mathcal{C} = 1$) for four-fermion, dipole, and Higgs-current operators. Current bounds (hatched bars), projected sensitivities (solid bars), and FCC-ee projections (shaded regions) are shown for $\tau \to \mu$ transitions.*

[Allwicher, Bartocci, Cornella, Leal, Paradisi, Scantamburlo, Sumensari, to appear.]

Message: LFV tau decays are more powerful than Z decays at FCC-ee!

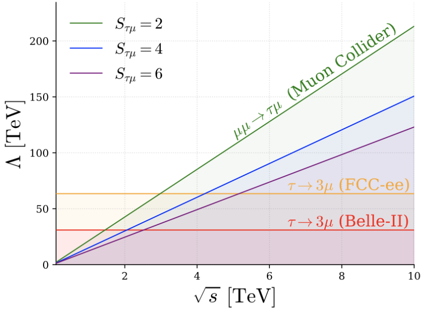


FIG. 7. Constraints on Λ associated with 4-fermion operators from Belle II, FCC-ee, and a Muon Collider. The signal significance satisfies $S = N_S/\sqrt{N_S + N_B}$, where N_S (N_B) denoting the numbers of signal (background) events.

Energy dependence of LFV signals

$$\Gamma(\tau \to 3e) \sim \frac{m_\tau^5}{\Lambda^4} \qquad \Gamma(Z \to \tau e) \sim \frac{m_Z^5}{\Lambda^4} \qquad \sigma(\mu\mu \to \tau e) \sim \frac{E_{CM}^2}{\Lambda^4}$$

Message: Energy helps accuracy!

[Allwicher, Bartocci, Cornella, Leal, Paradisi, Scantamburlo, Sumensari, to appear.]

LFV & the light new physics landscape [Opferkuch @ LNF, 2025]

Many Options, Much Fun?

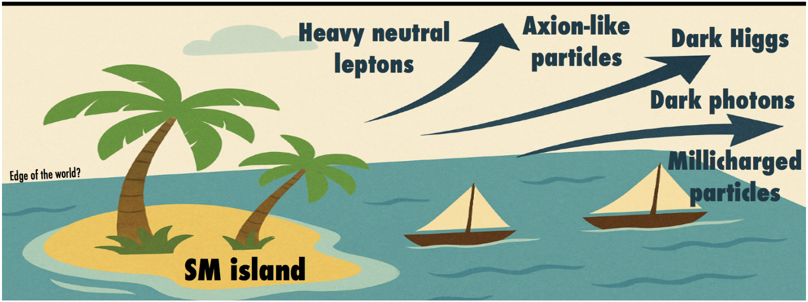
How do I make sense of the landscape of singlet options?

EFT: $\mathcal{L}_{\text{portal}} \supset \frac{c_n}{\Lambda^{\Delta_{\text{dark}}+\Delta_{\text{SM}}-4}} \mathcal{O}_{\text{dark}} \mathcal{O}_{\text{SM}}$

↳ Any gauge invariant combination of SM fields

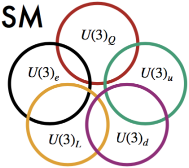
Lower dimensional portals
$\mathcal{O}_{\text{dark}} = \{\phi, N, V_\mu, \dots\}$
Scalar, Fermion, Vector

⟶ **Largest rates**



Flavoured Axion-like Portal

Accidental symmetries of the SM might be broken by light new particles feebly coupled to the SM



Example: Flavour dependent Peccei-Quinn charges

[Calibbi et al 1612.08040]
[Ema et al 1612.05492]

Flavour conserving + **Flavour violating**

$$\mathscr{L} \supset \sum_i \frac{\partial_\mu a}{2f_a} \bar{f}_i C^A_{ii} \gamma_\mu \gamma_5 f_i + \sum_{i \neq j} \frac{\partial_\mu a}{2f_a} \bar{f}_i \gamma^\mu \left(C^V_{ij} + C^A_{ij} \gamma_5 \right) f_j$$

[Feng et al hep-ph/9709411]

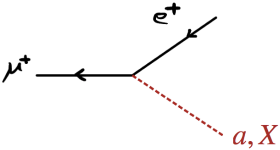
⚠ Hierarchy between **flavour-conserving** and **flavour-violating** depends on the UV theory

Flavour anarchy: $C^{A,V}_{ij}(\Lambda_{\rm UV}) \sim \mathcal{O}(1)$

Minimal favour violation: $C^{A,V}_{ij}(\Lambda_{\rm UV}) = 0$

What New Physics?

Lepton-flavour Violating

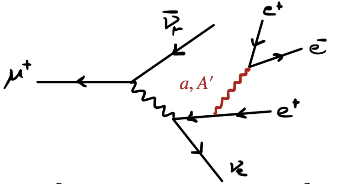


$$\mathscr{L}_{\text{FV}-\text{a}} = \frac{m_\mu}{2f_a}\frac{1}{|C_{e\mu}|}a\bar{\mu}(C^V_{e\mu} + C^A_{e\mu}\gamma_5)e$$

Bump hunt in electron momentum p_e

[Bayes et al (TWIST Collaboration) 1411.1770]
[Perrevoort et al (Mu3e Collaboration) 1812.00741]

Lepton-flavour Conserving



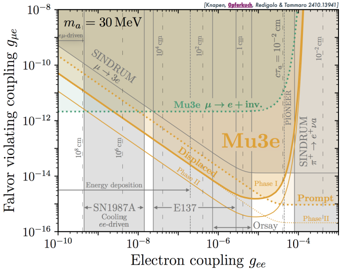
$$\mathscr{L}_0 \supset X\left[\bar{\mu}(g^\mu_S + g^\mu_P\gamma_5)\mu + \bar{e}(g^e_S + g^e_P\gamma_5)e\right]$$

$$\mathscr{L}_1 \supset X^\alpha\left[\bar{\mu}\gamma_\alpha(g^\mu_V + g^\mu_A\gamma_5)\mu + \bar{e}\gamma_\alpha(g^e_V + g^e_A\gamma_5)e\right]$$

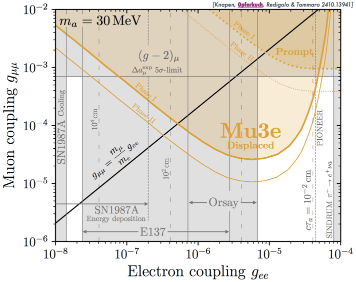
Bump hunt in e^+e^--pair invariant mass

[Echenard et al 1411.1770]
[Perrevoort et al (Mu3e Collaboration) 1812.00741]

Lepton Flavour Violating ALP

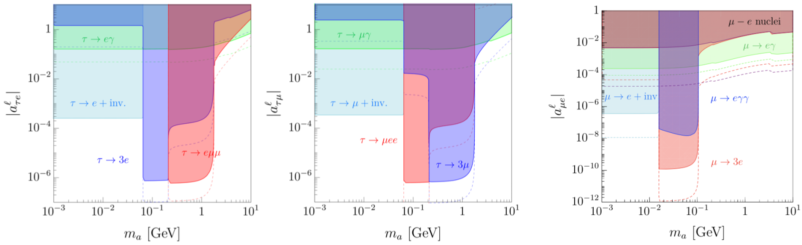


Lepton Flavour Conserving ALP

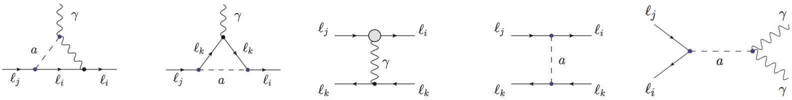


LFV from light NP: axionlike particles

$$\mathcal{L}_{\text{eff}}^{d\leq 5} = \frac{1}{2}(\partial_\mu a)(\partial^\mu a) - \frac{m_a^2 a^2}{2}$$
$$+ e^2\, c_{\gamma\gamma}\frac{a}{\Lambda}F_{\mu\nu}\widetilde{F}^{\mu\nu} + g_s^2\, c_{gg}\frac{a}{\Lambda}G_{\mu\nu}\widetilde{G}^{\mu\nu} - \frac{\partial_\mu a}{\Lambda}\sum_{f,i,j}\bar{f}_i\gamma^\mu(v_{ij}^f - a_{ij}^f\gamma_5)f_j$$



[Cornella, P.P. & Sumensari, '19.]



$$\mathcal{B}(\ell_i \to \ell_j\gamma\gamma) \approx \mathcal{B}(\ell_i \to \ell_j a) \times \mathcal{B}(a \to \gamma\gamma) \;\Rightarrow\; \mathrm{BR}(\ell_i \to \ell_j\ell_k\bar{\ell}_k) \not\approx \alpha \times \mathrm{BR}(\ell_i \to \ell_j\gamma)$$

Message: correlations among LFV signals discriminate heavy vs. light NP!

Conclusions and future prospects

- **Important questions in view of ongoing/future experiments are**:
  - What are the expected deviations from the SM predictions induced by TeV NP?
  - Which observables are not limited by theoretical uncertainties?
  - In which case we can expect a substantial improvement on the experimental side?
  - What will the measurements teach us if deviations from the SM are [not] seen?

- **(Personal) answers**:
  - We can expect any deviation from the SM expectations below the current bounds.
  - LFV processes and leptonic EDMs do not suffer from theoretical limitations and there are still excellent prospects for experimental improvements.
  - FCC-ee and a high-energy Muon Collider are complementary with each other and also with Belle II and LHCb to probe LFV.

Message: an exciting program is in progress at the Intensity Frontier which would greatly benefit from the next high-energy frontier program (FCC & MuC)!