



# Interacting orders in functional materials

Louis Ponet

#### **ISTITUTO ITALIANO DI TECNOLOGIA**

### Cr: charge and spin density wave, and ultrafast control

- Fermi surface nesting: High  $T_N$  spin density wave (311 K)
- Exchange striction  $\rightarrow$  charge density wave

$$
H = \sum_{\langle i,j \rangle} J(r_{i,j}) S_i S_j; \qquad r_{i,j} = |r_i - r_j|
$$
  

$$
J(r_{i,j}) = J(r_{i,j}^0) + \frac{dJ}{dr_{i,j}}\Big|_{r=r^0} \delta r
$$









### Cr: experimental setup

- Two optical pump pulses (40 fs width, 1-2 ps between pulses)
- x-ray probe monitors strain wave/CDW, created by SDW
- Pulses heat up SDW
- Electrostriction  $\rightarrow$  CDW amplitude oscillations





### Cr: experimental results

• 2 pulses: Constructive or destructive interference





### Cr: theory

- CDW  $\Leftrightarrow$  SDW, 1<sup>st</sup> order transition
- Two order parameters:  $L, y$

• 
$$
F = \frac{\alpha}{2} (T_L - T_c)L^2 + \frac{\beta_L}{4} L^4 - gL^2 y + \frac{\omega^2}{2} y^2 + \frac{\beta_y}{4} y^4
$$

- Two temperatures:  $T_b$ ,  $T_L$
- Pulses applied to  $T_{\rm L}$







### Cr: theoretical results





## Magnetism

- Interaction between spin order and orbital order in high spin-orbit coupled materials
- Heavy ions, e.g. Sr2IrO4, NaIrO3
- $\lambda$   $\boldsymbol{l} \cdot \boldsymbol{S}$ ;  $\lambda \sim Z^2$  $|\tilde{\uparrow}\rangle = sin(\theta) |0,\uparrow\rangle - cos(\theta) |+1,\downarrow\rangle$  $\left|\tilde{\downarrow}\right\rangle = sin(\theta) |0,\downarrow\rangle - cos(\theta) |-1,\uparrow\rangle$ •  $S \rightarrow \tilde{S}$

 $|\tilde{\uparrow}\rangle$  =  $|0,\uparrow\rangle$  +  $|1,\downarrow\rangle$ 





### Model

• Extended Heisenberg Hamiltonian

$$
H = \sum_{} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j \quad \longrightarrow \quad \mathcal{H}_{exch} = \sum_{i,j} \vec{m}_i \mathcal{J}_{ij} \vec{m}_j
$$

• Includes terms like Kitaev exchange



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## Numerical Calculation

• Based on Green's functions

 $\delta \vec{\phi_i}$ 

• Collinear, spin-only case (purely isotropic):

 $\mathcal{H}_0$ 

$$
\delta \vec{\mu}_i = \frac{1}{4\pi} \int_{-\infty}^{\delta \vec{\phi}_j} d\epsilon \, SpIm\left[\Delta_i G_{ij}^{\dagger} \Delta_j G_{ji}^{\dagger}\right]
$$

$$
J_{ij} = \frac{1}{4\pi} \int_{-\infty}^{E_f} d\epsilon \, SpIm\left[\Delta_i G_{ij}^{\dagger} \Delta_j G_{ji}^{\dagger}\right]
$$



KKR: A.I. Liechtenstein et al. Journal of Magnetism and Magnetic Materials (1987), *67*, 65–74. LDA+U: V. V. Mazurenko, *Phys. Rev. B, 71*(18), 1–8, *(2005)*, . LDA+U & Wannier90: D. M. Korotin, Phys. Rev. B (91), *224405*, 1–7 (2015).



## Numerical Calculation

- SOC case:
	- $-\sigma$  is not a good quantum number  $\rightarrow i$
	- $-$  Spins might be noncolinear  $\rightarrow$  local z-axis changes
	- $-$  Rotation of *j* rotates charge, changes hopping,  $G$ – Lot bigger space to rotate in  $\delta \mathcal{H}_i = 1/2i \delta \vec{\phi}_i [\mathcal{H}_i, \vec{j}_i]$
- Solution: rotate effective on-site  $B$  field





### Domain wall-induced elastic softening

- Interaction between ferroelectric polarization and strains
	- Apparent mechanical softening





### Model

- Interaction with Structural order
- Model: Ginzburg-Landau-Devonshire





## Model

- Interaction with Structural order
- Model: Ginzburg-Landau-Devonshire



Park, S. M. et al, (2018).*Nature Nanotechnology*, *13*(5), 366–370.







### Simulations

### • Finite element method

