



Interacting orders in functional materials

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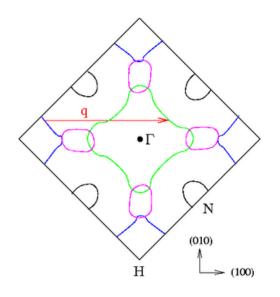
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Cr: charge and spin density wave, and ultrafast control

- Fermi surface nesting: High T_N spin density wave (311 K)
- Exchange striction \rightarrow charge density wave

$$H = \sum_{\langle i,j \rangle} J(r_{i,j}) S_i S_j; \qquad r_{i,j} = |\mathbf{r}_i - \mathbf{r}_j|$$
$$J(r_{i,j}) = J(r_{i,j}^0) + \frac{dJ}{dr_{i,j}} \bigg|_{r=r^0} \delta r$$



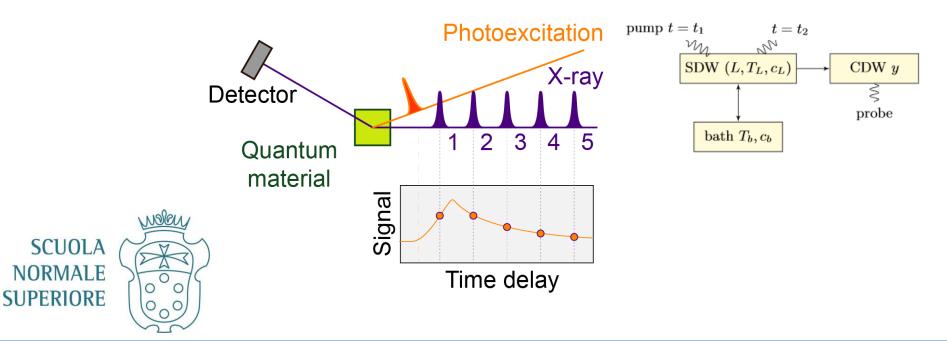






Cr: experimental setup

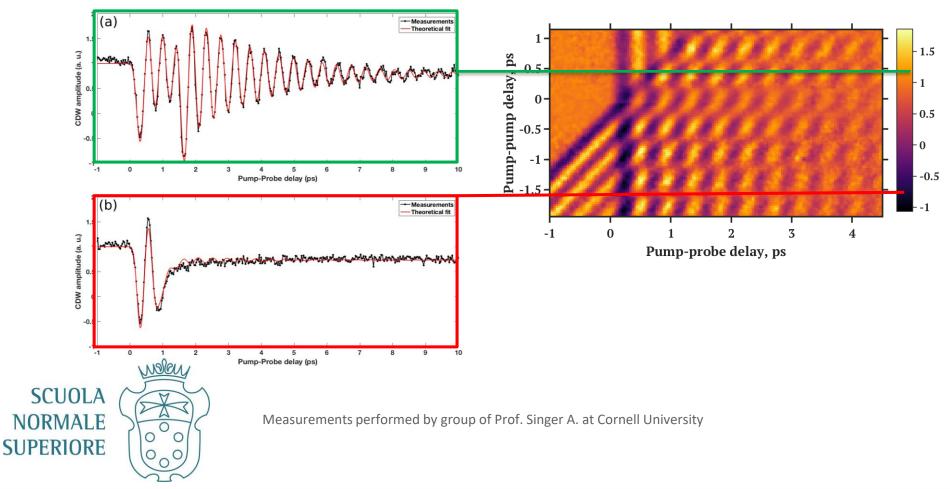
- Two optical pump pulses (40 fs width, 1-2 ps between pulses)
- x-ray probe monitors strain wave/CDW, created by SDW
- Pulses heat up SDW
- Electrostriction \rightarrow CDW amplitude oscillations





Cr: experimental results

• 2 pulses: Constructive or destructive interference



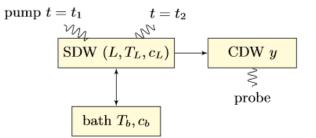


Cr: theory

- CDW ⇔ SDW, 1st order transition
- Two order parameters: *L*, *y*

•
$$F = \frac{\alpha}{2} (T_L - T_C) L^2 + \frac{\beta_L}{4} L^4 - g L^2 y + \frac{\omega^2}{2} y^2 + \frac{\beta_y}{4} y^4$$

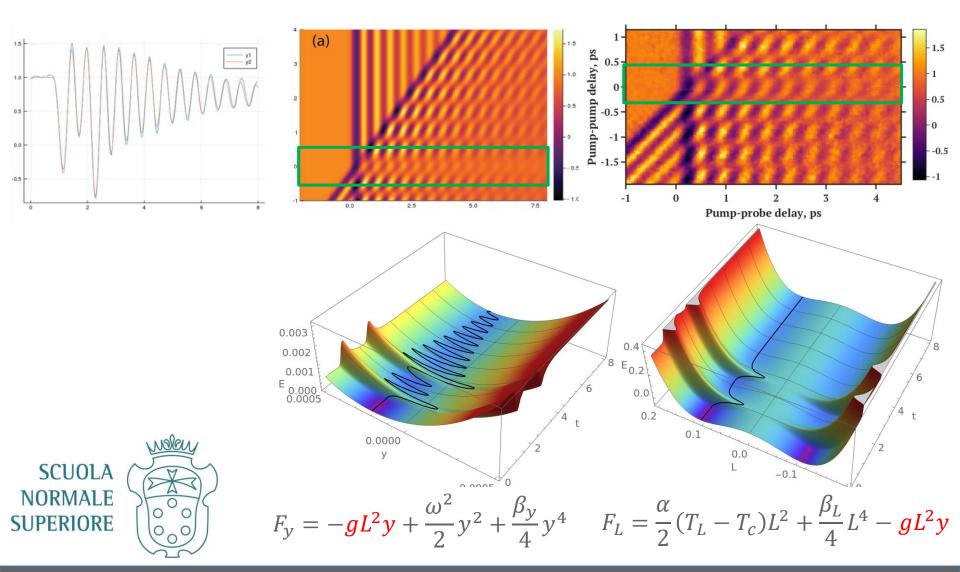
- Two temperatures: T_b , T_L
- Pulses applied to $T_{\rm L}$







Cr: theoretical results





Magnetism

- Interaction between spin order and orbital order in high spin-orbit coupled materials
- Heavy ions, e.g. Sr2IrO4, NaIrO3
- $\lambda l \cdot S; \lambda \sim Z^2$ $|\tilde{\uparrow}\rangle = sin(\theta) |0, \uparrow\rangle cos(\theta) |+1, \downarrow\rangle$ • $S \rightarrow \tilde{S}$ $|\tilde{\downarrow}\rangle = sin(\theta) |0, \downarrow\rangle - cos(\theta) |-1, \uparrow\rangle$

 $\left|\tilde{\uparrow}\right\rangle$ = $\left|0,\uparrow\right\rangle$ + $\left|1,\downarrow\right\rangle$





Model

• Extended Heisenberg Hamiltonian

$$H = \sum_{\langle ij \rangle} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j \qquad \longrightarrow \qquad \mathcal{H}_{exch} = \sum_{i,j} \vec{m}_i \mathcal{J}_{ij} \vec{m}_j$$

• Includes terms like Kitaev exchange



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Numerical Calculation

Based on Green's functions

 $\delta ec{\phi_i}$

Collinear, spin-only case (purely isotropic):

 \mathcal{H}_0

$$\delta \vec{\phi}_{j} \qquad \delta \mathcal{H}_{i} = 1/2i\delta \vec{\phi}_{i}[\mathcal{H}_{i}, \vec{\sigma}_{i}] \\ \Delta_{i} = \mathcal{H}_{i}^{\uparrow\uparrow} - \mathcal{H}_{i}^{\downarrow\downarrow} \\ \delta \mathcal{H}_{i} = \frac{\Delta_{i}}{2} \left(\delta \phi_{i}^{y} \sigma_{i}^{x} - \delta \phi_{i}^{x} \sigma_{i}^{y}\right) \\ J_{ij} = \frac{1}{4\pi} \int_{-\infty}^{E_{f}} d\epsilon \, Sp \, Im \left[\Delta_{i} G_{ij}^{\downarrow} \Delta_{j} G_{ji}^{\uparrow}\right]$$



KKR: A.I. Liechtenstein et al. Journal of Magnetism and Magnetic Materials (1987), 67, 65–74. LDA+U: V. V. Mazurenko, *Phys. Rev. B, 71*(18), 1–8, (2005), . LDA+U & Wannier90: D. M. Korotin, Phys. Rev. B (91), 224405, 1–7 (2015).



Numerical Calculation

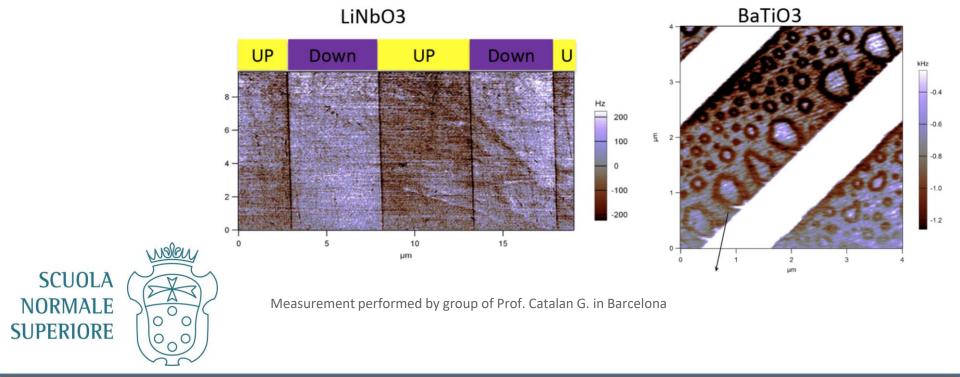
- SOC case:
 - $-\sigma$ is not a good quantum number $\rightarrow j$
 - Spins might be noncolinear \rightarrow local z-axis changes
 - Rotation of j rotates charge, changes hopping, G
 - Lot bigger space to rotate in $\delta \mathcal{H}_i = 1/2i\delta \vec{\phi}_i [\mathcal{H}_i, \vec{j}_i]$
- Solution: rotate effective on-site **B** field





Domain wall-induced elastic softening

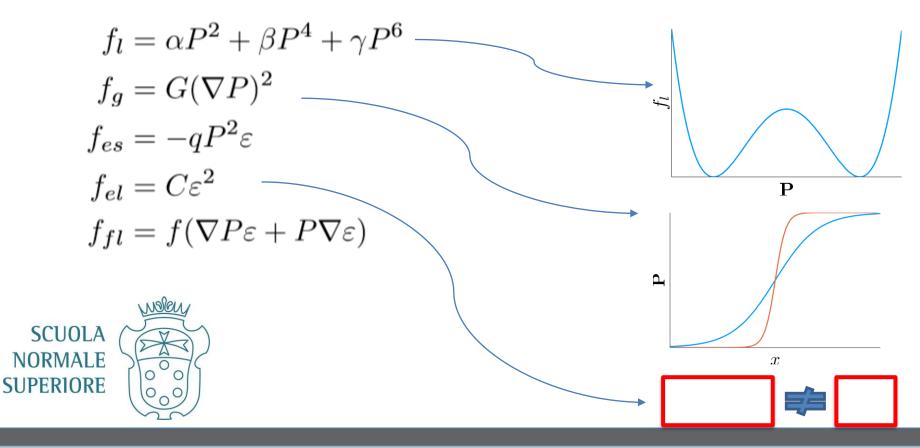
- Interaction between ferroelectric polarization and strains
 - Apparent mechanical softening





Model

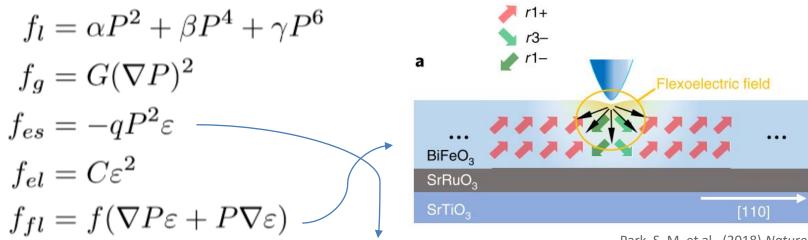
- Interaction with Structural order
- Model: Ginzburg-Landau-Devonshire



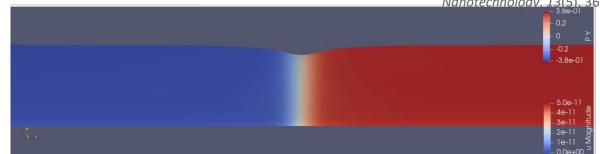


Model

- Interaction with Structural order
- Model: Ginzburg-Landau-Devonshire



Park, S. M. et al, (2018).*Nature Nanotechnoloav*, *13*(5), 366–370.







Simulations

• Finite element method

