

HOLOGRAPHY AND KKLT

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GRAVITY, STRING AND SUPERSYMMETRY BREAKING

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CHALLENGES FOR STRING PHENOMENOLOGY

Issues from compactification:

1. Deformations of compactification space

→ lots of massless scalar fields (moduli)
(Dine-Seiberg-like control issues)

2. Natural: AdS or Minkowski ($\Lambda \leq 0$)

de Sitter vacua ($\Lambda > 0$): extremely difficult

(Conjecture: no (long lived) de Sitter in Quantum Gravity [Vafa et al. '18, '19])

CONTROLLED STRING THEORY DE SITTER VACUA

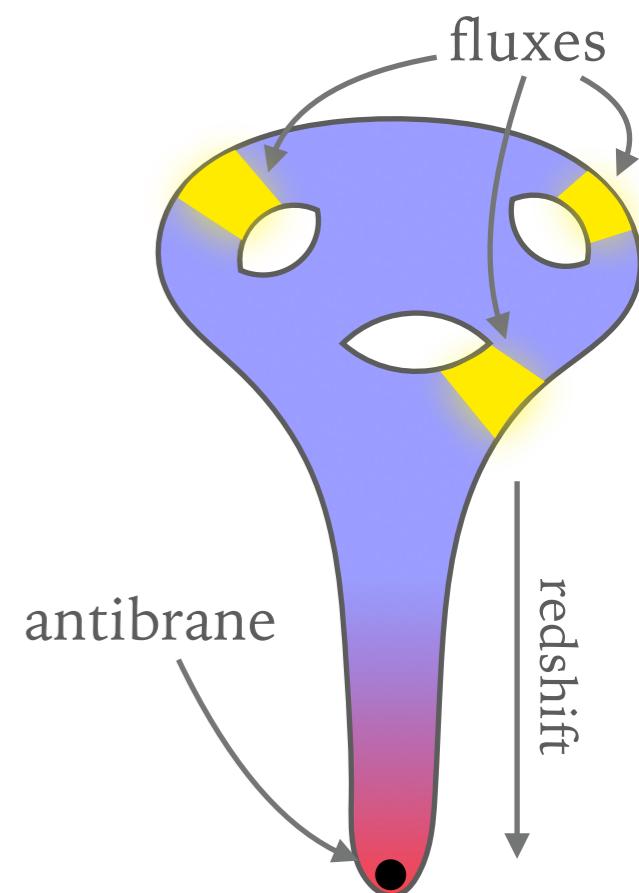
KKLT scenario: two step procedure:

1. stabilize moduli with fluxes + non-pert. effects

→ SUSY vacuum with
 $\Lambda < 0$ (AdS)

2. raise cosmological constant above zero
 (perturb by antibrane in warped throat)

→ broken SUSY and
 $\Lambda > 0$



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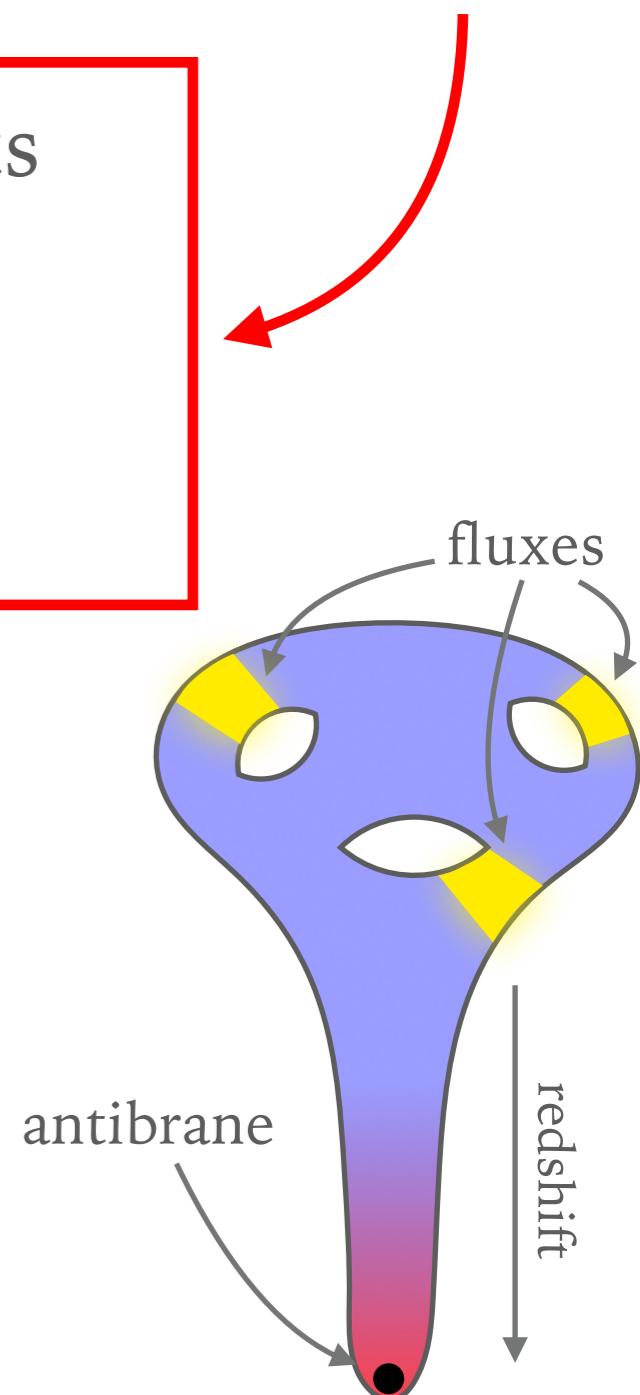
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This Talk!



REVIEW:

KKLT ADS VACUA

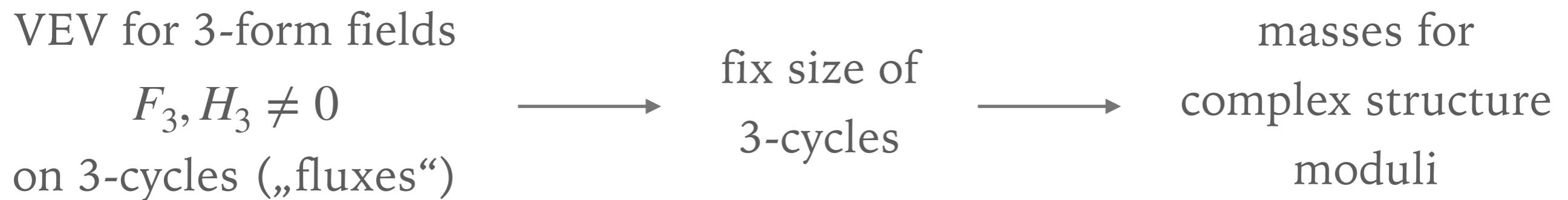
COMPACTIFICATION ON CALABI-YAU MANIFOLDS

- IIB on Calabi-Yau (orientifolds):

$$M_{10} = M_4 \times_w CY_3$$

- two classes of moduli:

- $h^{1,1}$ Kähler moduli (volumes of 2 or 4-cycles)
- $h^{2,1}$ complex structure moduli (volumes of 3-cycles)



MODULI STABILIZATION

- tree-level $\mathcal{N} = 1$ superpotential:

$$W = \int G_3 \wedge \Omega$$
$$(G_3 = F_3 - \tau H_3)$$

Does not depend on
Kähler moduli!

- F-term conditions

complex structure

$$D_i W = 0$$

Kähler

$$D_\alpha W \sim W$$

+ axio-dilaton:

$$\Rightarrow \star G_3 = i G_3$$

moduli:

\Rightarrow SUSY broken
unless $W = 0$

- Resolution: non-perturbative quantum corrections:

$$W = \int G_3 \wedge \Omega + \sum_{\mathbf{k}} \mathcal{A}_{\mathbf{k}}(z^i, G_3) e^{-2\pi k^\alpha T_\alpha}$$

Kähler moduli

→ W depends on Kähler moduli!

KKLT: AdS VACUUM

[Kachru, Kallosh, Linde, Trivedi '03]

► Full scalar potential:

$$V = e^K \left(g^{a\bar{b}} D_a W \bar{D}_{\bar{b}} \bar{W} - 3 |W|^2 \right)$$

► Supersymmetry conditions:

(for all moduli)

$$D_a W = 0$$

Notice:

*Existence of classical flux vacuum
does not guarantee existence of
KKLT AdS vacuum!*

► Supersymmetric AdS vacuum:

$$\Lambda_{AdS} = \langle V \rangle = -3 \left(e^K |W|^2 \right) \Big|_{D_a W = 0}$$

*control over
instanton expansion:*

$$\left| e^{-2\pi k^\alpha T_\alpha} \right| \ll 1$$



*controlled uplift to dS
only possible if*

$$\Lambda_{AdS} \ll 1$$

M-THEORY / F-THEORY

- M-theory on **Calabi-Yau 4-fold** (from 11D to 3D)

$h^{3,1}$ complex str. moduli
(volumes of 4-cycles)



stabilized by
 $G_4 \neq 0$
(4-form flux)

- Classical superpotential [Gukov, Vafa, Witten '99]:

$$W \sim \int_{CY_4} G_4 \wedge \Omega$$

- Non perturbative corrections:

$$W = \int G_4 \wedge \Omega + \sum_{\mathbf{k}} \mathcal{A}_{\mathbf{k}}(z^i, G_4) e^{-2\pi k^\alpha T_\alpha}$$

M-THEORY / F-THEORY

- classical flux vacua:

$$D_i W = \int G_4 \wedge \chi_i = 0 \quad \Rightarrow \quad \begin{aligned} G_4 &= \star G_4 \\ G_4 &\in H^{4,0} \oplus H_+^{2,2} \oplus H^{0,4} \end{aligned}$$

- KKLT like AdS vacua:

balance ↪

- $G_4^{4,0} \neq 0$
- non-pert. corrections

$$\langle V \rangle = -4 \left(e^K |W|^2 \right) \Big|_{D_a W=0} < 0$$

- Tadpole cancellation [Becker, Becker '96; Sethi, Vafa, Witten '96]:

$$N_{M2} + \frac{1}{2} \int_{CY_4} G_4 \wedge G_4 = \frac{\chi(CY_4)}{24}$$

$\chi(CY_4)$: Euler number of CY 4-fold
(F-theory: encodes IIB D7-charges)

HOLOGRAPHY AND KKLT

HOLOGRAPHIC PERSPECTIVE

- first step of KKLT requires
 - AdS vacuum with
$$\Lambda_{AdS} \ll 1$$
 - (consistency of EFT: $\Lambda_{AdS} \ll m_{KK}^2 \rightarrow$ scale separation)
- Is there a holographically dual CFT?
 - AdS/CFT duality:
- Dual CFT must have $c \gg 1$

HOLOGRAPHIC PERSPECTIVE

- first step of KKLT requires

AdS vacuum with

$$\Lambda_{AdS} \ll 1$$

(consistency of EFT: $\Lambda_{AdS} \ll m_{KK}^2 \rightarrow$ scale separation)

- Is there a holographically dual CFT?

AdS/CFT duality:

$$|\Lambda_{AdS}| \sim \frac{1}{c}$$

central charge
(# degrees of freedom of CFT)

- Dual CFT must have $c \gg 1$

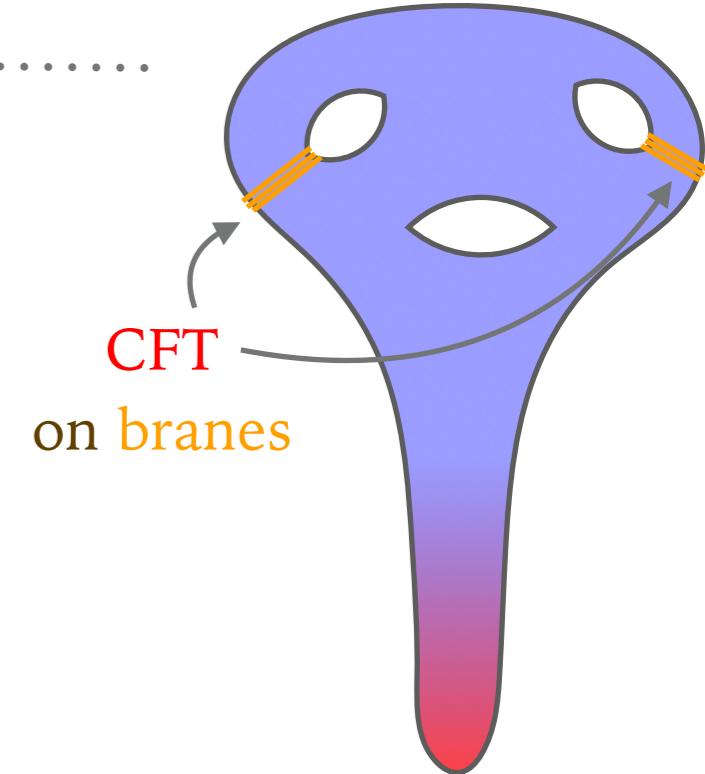
HOLOGRAPHICALLY DUAL BRANE THEORIES

Conventional Holography:

*Dualize Fluxes
into Branes*



*brane world volume
theory flows to CFT*



► IIB with G_3 flux (KKLT):

D5 & NS5
branes



on special
Lagrangian 3-cycles
of CY_3



SCFT in
 $d = 3, \mathcal{N} = 1$

► M-theory with G_4 flux (F-theory):

M5
branes



on special
Lagrangian 4-cycles
of CY_4



SCFT in
 $d = 2, \mathcal{N} = (1,1)$

FLUXES AND DUAL BRANES

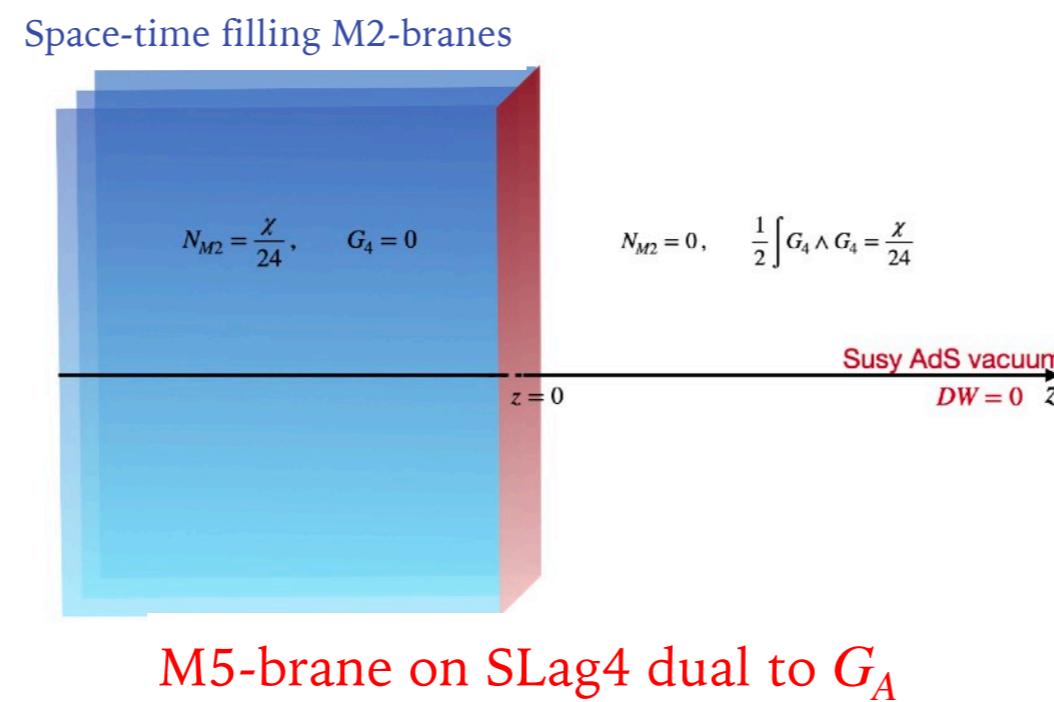
- Consider two different flux vacua with $G_4 = G_A$ and $G_4 = G_B$ on the same CY (+ different numbers of M2-branes).
- M5 brane on 4-cycle C dual to $[G_A - G_B]$ [Gukov, Vafa, Witten '99]
 - Domain Wall
 - interpolating between the two vacua

- Extremal case: $G_B = 0$
 - [SL, Vafa, Wiesner, Xu '22]
 - [Bena, Li, SL '24]

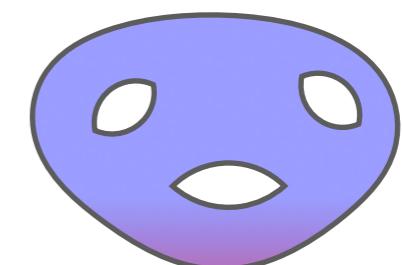
$W_{GVW} = 0$
*existence of minimum at
finite CY-volume unclear*



??



KKLT AdS-vacuum



$|\Lambda| \ll 1$

TOPOLOGICAL COUNT OF DEGREES OF FREEDOM

COUNTING DEGREES OF FREEDOM

[SL, Vafa, Wiesner, Xu '22]

- Brane world volume theory



- c-theorem:
- AdS/CFT:

$$c_{UV} > c_{IR}$$

- Goal: Count degrees of freedom in UV



COUNTING DEGREES OF FREEDOM

[SL, Vafa, Wiesner, Xu '22]

- Brane world volume theory



- **c-theorem:**
 $c_{UV} > c_{IR}$
- **AdS/CFT:**
 $c_{IR} \sim |\Lambda_{AdS}|^{-1}$
- Goal: Count degrees of freedom in UV



COUNTING DEGREES OF FREEDOM (2)

[SL, Vafa, Wiesner, Xu '22]

- M5 brane on sLag cycle L_4 : **world volume theory**:

reduction of six-dim.

$\mathcal{N} = (2,0)$ theory on L_4 \longrightarrow two-dimensional
 $\mathcal{N} = (1,1)$ theory

count moduli of L_4 \Rightarrow $c_{UV} = \frac{3}{2}(2 + b_2^+ + b_2^- + 2b_1) \lesssim \alpha L_4 \cdot L_4$

self-intersection of L_4 \curvearrowright

b_n : Betti numbers of L_4 \curvearrowleft

- Relation with **Tadpole Condition**:

$$\frac{1}{2}L_4 \cdot L_4 = \frac{1}{2} \int_{CY_4} G_4 \wedge G_4 \leq \frac{\chi(CY_4)}{24} \Rightarrow$$

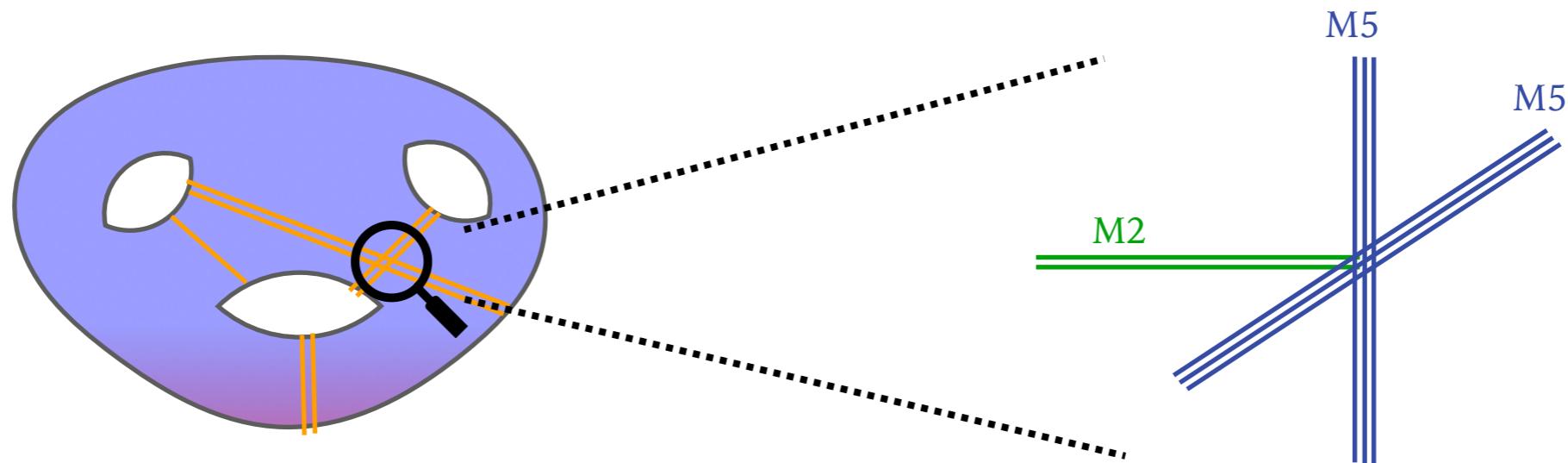
$$c_{IR} \lesssim \beta \chi(CY_4)$$

DEGREES OF FREEDOM IN SUPERGRAVITY

COUNTING OF DEGREES OF FREEDOM IN SUPERGRAVITY

[Bena, Li, SL '24]

- Most degrees of freedom are localized at brane intersections:

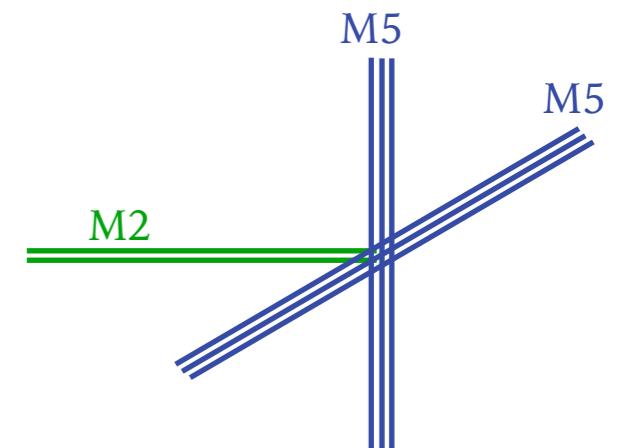


- Local description: Intersecting branes in flat space
- Goal: Find supergravity solution and count degrees of freedom holographically

INTERSECTING BRANES IN SUPERGRAVITY

- Local brane setup:

	t	y	z	1	2	3	4	5	6	7	8
M5	x	x		x	x	x	x				
M5	x	x						x	x	x	x
M2	x	x	x								



- Supergravity solution in the near horizon limit:

$$ds_{11}^2 = \underbrace{\left(\frac{u}{l_{\text{AdS}}} \right) (-dt^2 + dy^2)}_{\text{AdS}_3} + \underbrace{\left(\frac{l_{\text{AdS}}}{u^2} \right)^2 du^2}_{S^3 \times S^3} + \underbrace{(n_T n_F)^{\frac{1}{3}} (d\Omega^2 + d\Omega'^2)}_{\mathbb{R}^2} + d\lambda^2 + \underbrace{\left(\frac{n_F^2}{n_T} \right)^{\frac{2}{3}} dz^2}_{\mathbb{R}^2}$$

with $l_{\text{AdS}} = \frac{1}{\sqrt{2}} (n_T n_F)^{\frac{1}{6}}$

[de Boer, Pasquinucci, Skenderis '99]

HOLOGRAPHIC CENTRAL CHARGE

[Bena, Li, SL '24]

- Normalization of brane charges:

compactify formally

$$\mathbb{R}^2 \rightarrow S_\lambda \times S_z$$



with radii: R_λ, R_z

(later: $R_\lambda, R_z \rightarrow \infty$)

$$N_5 = \frac{1}{(2\pi)^3} \int_{S_z \times S^3} G_4 = R_z n_F$$

$$N_2 = \frac{1}{(2\pi)^6} \int_{S_\lambda \times S^3 \times S^3} (\star G_4 - \frac{1}{2} C_3 \wedge G_4) = \frac{\sqrt{2}}{4\pi} R_\lambda (n_T n_F)^{\frac{5}{6}}$$

- Central charge of dual CFT_2 :

$$l_{AdS_3} = \frac{1}{\sqrt{2}} (n_T n_F)^{\frac{1}{6}}$$

$$c = \frac{3}{2} \frac{l_{AdS_3}}{G_N^{(3)}} = 3 N_2 N_5 \sim N_5^3$$

$$\begin{aligned} \frac{1}{G_N^{(3)}} &= \text{vol}(S^3 \times S^3 \times S_\lambda \times S_z) \\ &\sim R_z R_\lambda (n_T n_F)^{\frac{2}{3}} n_T \end{aligned}$$

Tadpole cancellation:

$$N_2 = \frac{\chi(CY_4)}{24} \sim N_5^2$$

- Generic situation: branes intersect at many points:

each intersection:

$$c \sim \mathcal{O}(1)$$

\times #intersections: N_5^2



$$c \sim N_5^2 \sim \chi(CY_4)$$

IIB: INTERSECTING D5 / NS5 BRANES ON CY3-ORIENTIFOLDS

- Local brane setup:

	t	x	y	z	1	2	3	4	5	6
D5	x	x	x		x	x	x			
NS5	x	x	x					x	x	x
D3	x	x	x	x						

- Near brane geometry:
 - [D'Hoker, Estes, Gutperle '07 (2x)]
 - [Aharony, Berkooz, Shamir '11]
 - [Assel, Bachas, Estes, Gomis '11]

$$AdS_4 \times_w S^2 \times S^2 \times \Sigma_2$$

Riemann surface

- Scaling of central charge:

$$c \sim \frac{l_{AdS_4}}{G_N^{(4)}} \sim (N_{D5} N_{NS5})^2$$

CONSEQUENCES FOR KKLT

CONSEQUENCES

[SL, Vafa, Wiesner, Xu '22]

- central charge bounded by tadpole condition:

$$\Lambda_{AdS} \sim \frac{1}{c} \gtrsim \frac{1}{\chi(CY_4)}$$

Euler number:

$$\chi(CY_4) = 6(8 + h^{1,1} + h^{3,1} - h^{2,1})$$

(see also “tadpole conjecture” [Bena, Blåbäck, Graña, SL '20])

- effective cutoff in gravitational theories with N light d.o.f.

[Dvali '07]

species scale: $\Lambda_{\text{species}} \sim \frac{1}{N}$

Here:

$$N \sim \chi(CY_4)$$

$$\Rightarrow \quad \Lambda_{AdS} \gtrsim \Lambda_{\text{species}}$$

AdS scale is above cutoff of the EFT!
→ no sensible EFT description!

CONCLUSIONS

- Holographic dual to KKLT AdS vacua:

*world-volume SCFT of
M5 (M-theory) / D5-NS5 (IIB) branes*

- Holography:

*counting d.o.f. of
M5-brane CFT:*

$$|\Lambda_{AdS}| > \left(N_{\text{moduli}}\right)^{-1}$$

- comparison with species scale:

$$\Lambda_{AdS} \gtrsim \Lambda_{EFT}$$



*no scale separated
KKLT like AdS vacua*

THANK YOU!