Modern Cosmology: Facts, Misconceptions and Myths

V. Mukhanov

ASC, LMU, München

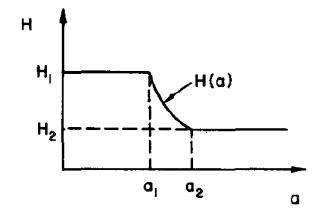
In desperation I asked Fermi whether he was not impressed by the agreement between our calculated numbers and his measured numbers. He replied, "How many arbitrary parameters did you use for your calculations?" I thought for a moment about our cut-off procedures and said, "Four." He said, "I remember my friend Johnny von Neumann used to say, with four parameters I can fit an elephant, and with five I can make him wiggle his trunk." Freeman Dyson, 1953

THEORY OF COSMOLOGICAL PERTURBATIONS

PHYSICS REPORTS (Review Section of Physics Letters) 215, Nos. 5 & 6 (1992) 203-333. North-Holland

V.F. MUKHANOV^{a,b}, H.A. FELDMAN^c and R.H. BRANDENBERGER^a

Par	t III.	Extensions	296
16.	Intro	duction	296
17.	Micro	owave background anisotropies	297
18.	Gravitational waves		301
	18.1.	Quantization	302
	18.2.	Observables	305
	18.3.	Spectrum of gravitational waves in de Sitter space	306
	18.4.	Spectrum of gravitational waves in the inflationary	
		universe	307
	18.5.	Spectrum of gravitational waves in double inflation	
		models	310
19.	Entropy perturbations		313
	19.1.	General remarks	313
	19.2.	A model for entropy perturbations	314
	19.3.	Evolution of the homogeneous field	314
	19.4.	Perturbations	316
	19.5.	Mountain and valley spectra	319
	19.6.	Suppression of long-wavelength perturbations	320
		Modulation of the spectrum in double inflation	
		models	321



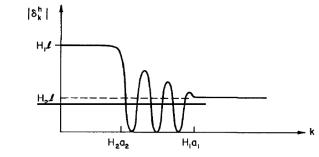


Fig. 18.7. Power spectrum $\delta_h = h$ of gravitational radiation in the double inflation model of (18.40) with the ratio of Hubble constants $H_1/H_2 = 10$.

Note that the effect discussed in this section can arise for both adiabatic and entropy perturbations. it is possible to obtain a suppression of the long-wavelength part of cosmological perturbation

$$S = \int \left[\frac{1}{2} \chi_{;\mu} \chi^{;\mu} - V(\chi) + \frac{1}{2} \varphi_{;\mu} \varphi^{;\mu} - \frac{1}{2} m_0^2 \varphi^2 - V_{\rm I}(\chi,\varphi,\ldots) \right] \sqrt{-g} \, \mathrm{d}^4 x \, ,$$

nontrivial spectra with mountains and valleys can also be obtained It is also possible to generate non-Gaussian fluctuations However, this procedure is extremely unappealing since it implies a complete loss of predictability.

MFB, Physics Report, 1992

Inflation is not a unique theory, but rather a class of models based on similar principles.

WRONG!

Inflation is THE theory only when it is understood as the stage of unbroken accelerated expansion due to the same ingridient which is responsible for quantum fluctuations.

Otherwise it is rubbish without any predictions!!!

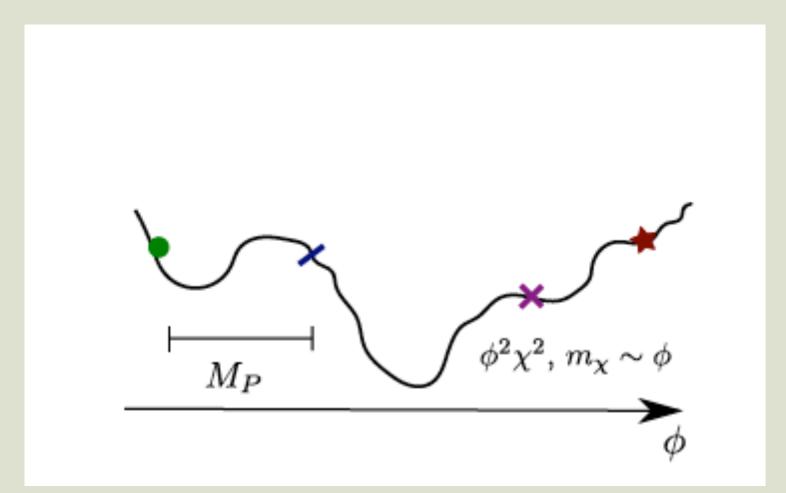
In this case it is unbeatable as predictive theory because it allows us to calculated the effect of amplification of quantum fluctuations in completely controlable weak coupling regimes

while most alternatives cannot even compete with "rubbish inflation" in a sense of controlable reproduction of outcome for quantum fluctuations

COSMOLOGY - Theology = $\exp(Ht)$ during at least 70 H^{-1} , but less than $10^6 H^{-1} \rightarrow$ no any problems with predictions, which could falsify the theory in Popper's sense The only purpose of inflationary models relevant for observation is a maping $V(\varphi)$ to $p \approx -\varepsilon$

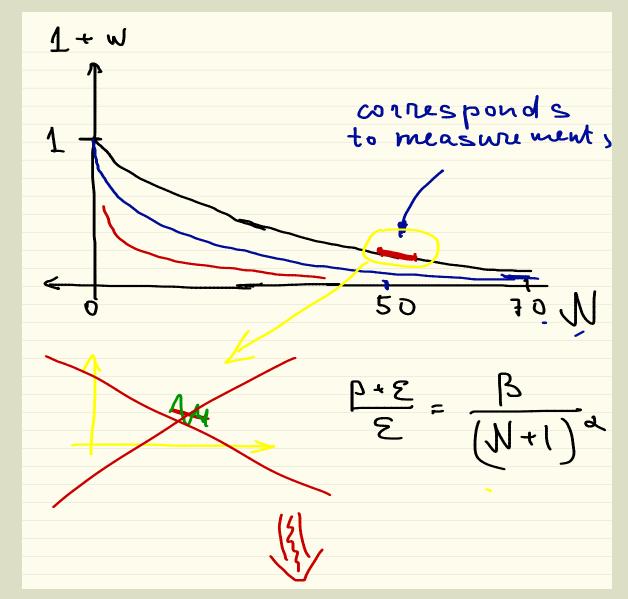
and this maping happened to be not crucial for robust predictions but important only for excluding definite potentials $V(\varphi)$, which anyway we will never be able to verify in any other independent experiments

$$V(\tau,\theta) = \frac{12W_0^2\xi}{(4\mathcal{V}_m - \xi)(2\mathcal{V}_m + \xi)^2} + \frac{D_1 + 12e^{-2a_2\tau}\xi A_2^2}{(4\mathcal{V}_m - \xi)(2\mathcal{V}_m + \xi)^2} + \frac{D_2 + \frac{16(a_2A_2)^2}{3a\lambda_2}\sqrt{\tau}e^{-2a_2\tau}}{(2\mathcal{V}_m + \xi)}$$
(25)
+
$$\frac{D_3 + 32e^{-2a_2\tau}a_2A_2^2\tau(1 + a_2\tau)}{(4\mathcal{V}_m - \xi)(2\mathcal{V}_m + \xi)} + \frac{D_4 + 8W_0A_2e^{-a_2\tau}\cos(a_2\theta)}{(4\mathcal{V}_m - \xi)(2\mathcal{V}_m + \xi)} \left(\frac{3\xi}{(2\mathcal{V}_m + \xi)} + 4a_2\tau\right) + \frac{\beta}{\mathcal{V}_m^2}.$$



What is relevant for predictions? $-\varepsilon$ energy density -p pressure $1 + w \equiv \frac{\mathcal{E} + p}{\ll} \ll 1$ during last 70 e-folds ($a = a_f \cdot e^{-N}$) *a*) $1 + w \ll 1$ for $N \gg 1$ b) $1 + w \approx O(1)$ for $N \simeq O(1)$ c) 1+w is a smooth function of N a) $1 + w \ll 1$ for $N \gg 1$

- b) $1 + w \approx O(1)$ for $N \simeq O(1)$
- c) 1+w is a smooth function of N



PREDICTIONS

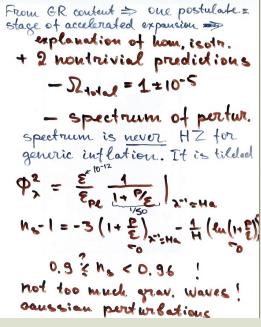
("smoking guns"-nonconfirming any of them would falsify THE theory)

- flat universe
- adiabatic perturbations
- small non-gaussianity ($f_{NL} \sim O(1)$)
- red-tilted spectrum

$$\Phi^2 \propto \lambda^{1-n_s}$$
$$1 - n_s = 3(1 + w) - \frac{d\ln(1 + w)}{dN}$$

of the faint ripples that we detect in the cosmic microwave background (CMB). First, the ripples should be nearly scale-invariant, meaning that they have nearly the same intensity at

The theory always predicts red-tilted spectrum



Cambridge, 2000

[1]. Contrary to an erroneous belief inflation does not predict the scale-invariant, Harrison-Zel'dovich spectrum. The spectral index should be in the range of $0.92 < n_s < 0.97$. The physical

V. Mukhanov, CMB, Quantum Fluctuations and the Predictive Power of Inflation, *arXiv* : *astro* – *ph*/0303077 (2003) Red-tilted log spectrum (MC, H, 1981-1982) \rightarrow

$$n_{\rm S} = 1 - \frac{A}{\ln(B\lambda_{\rm gal} / \lambda_{\rm CMB})},$$

where A > 1,5 and $B \simeq 1-100$ depending on $50 < N < 55 \rightarrow$

 $n_{\rm s} < 0.97$

irrespective of any particular model!

L.P. 9/6/2003:

We are writing a proposal to get money to do our small angular scale CMB experiment. If I say that simple models of inflation require $n_s=0.95+/-0.03$ (95\% cl) is it correct?

I'm especially interested in the error. Specifically, if n_s=0.99 would you throw in the towel on inflation?

V.M. 9/8/2003

The "robust" estimate for spectral index for inflation is $0.92 < n_s < 0.97$. The upper bound is more robust than lower. The physical reason for the deviation of spectrum from the flat one is the nessesity to finish inflation.... If you find $n_s=0.99 + -0.01$ (3 sigma) I would throw in the towel on inflation. Further predictions ("non-smoking guns"):

- Primordal gravitational waves
- Nongaussianities due to nonlinearity of
 Einstein equation (3,4,...points correlaton functions)

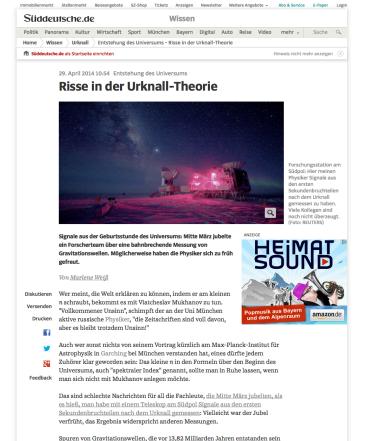
There must be primordial gravi tatonal waves $r \equiv \frac{T}{S} = 24 \cdot (1 + p / \varepsilon) = \frac{\beta}{N^{\alpha}}$ No a priori low bound on their ampltude!

A if no is measured fixed by amplo ns=0.96 upper limit (2-0.2) linit 2~0,004 Eend~10⁻¹² Epe

- However, keeping mind theoretical and experimental uncertainty, n_s within 2-sigma can be equal to 0.95.
- In this case the lower bound on r becomes 0.0006
- (unrealistic from the point of view of future measurements)

- Thus, detection of the primordial gravitational field will provide us an extra confirmation that quantum fluctuations were amplified on the stage of accelerated expansion.
- Failing to detect them at the level 0.04 would not have any implications and in no way can be considered as a prove of alternative for amplification of quantum fluctuation

I heory is right Plank is right BICEP2 is night $T+P \vee T+B \vee$ P+B V but TTTB Therefore P+B > catastrophy tor theory



Who thinks he can explain the world by screwing the small n, it gets to do with Viatcheslav Mukhanov. "Perfect nonsense," complains the active at the University of Munich Russian physicist, "the magazines are full of it, but still it remains nonsense!"

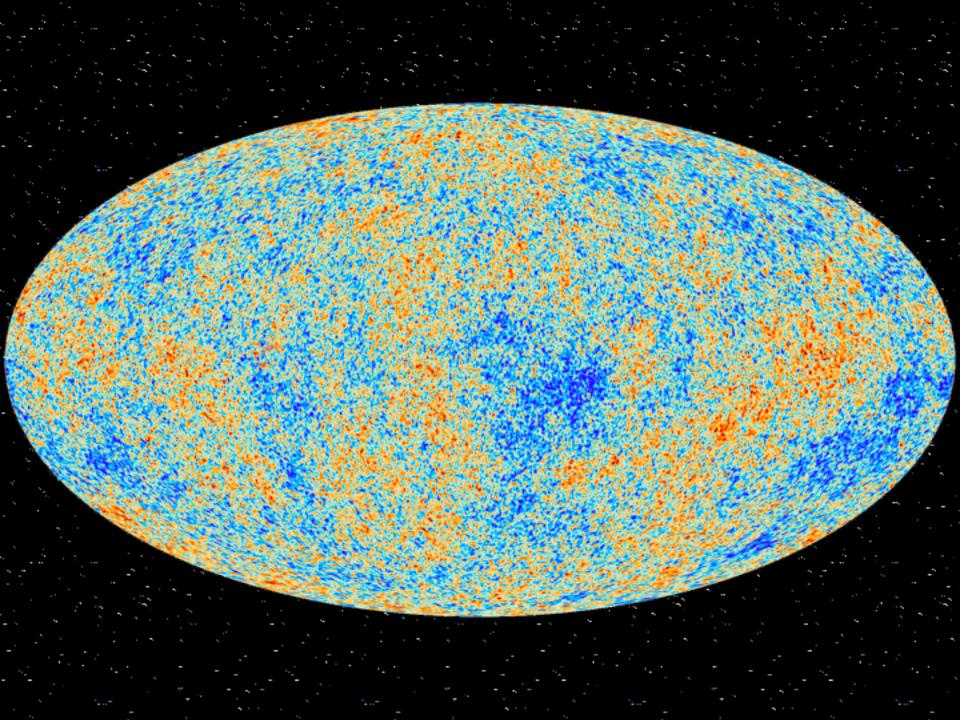
Non-gaussianities

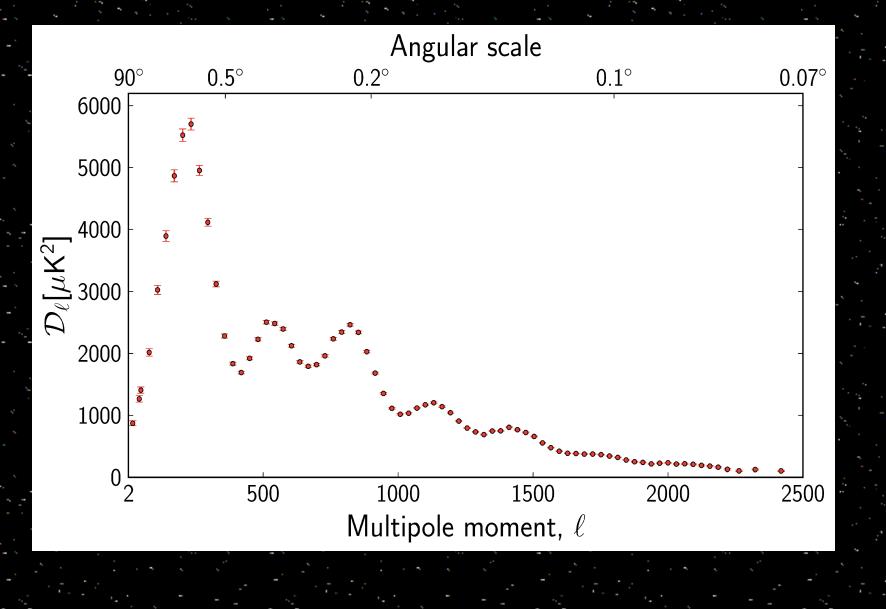
$$\Phi = \Phi_g + f_{NL} \Phi_g^2$$

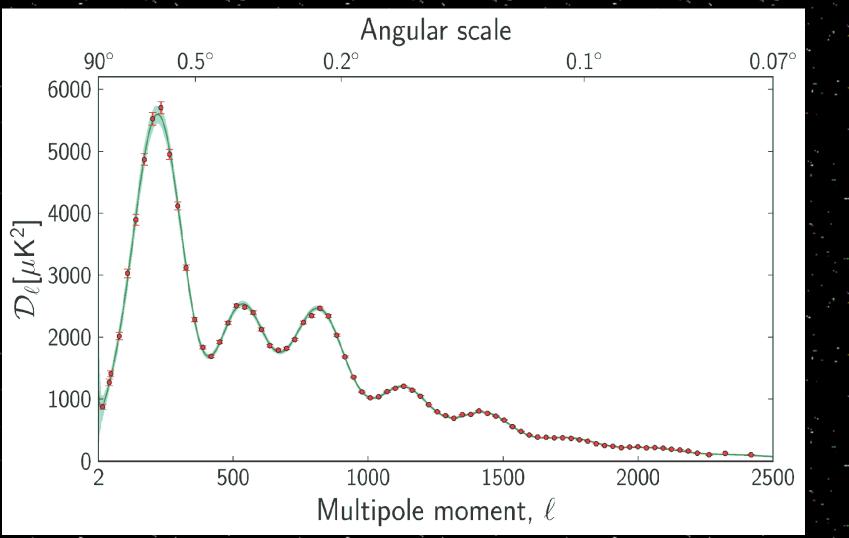
 $f_{NL} \simeq 0.04$ from inflation and $f_{NL} \simeq 2 - 4$ from subsequent evoluton of perturbations

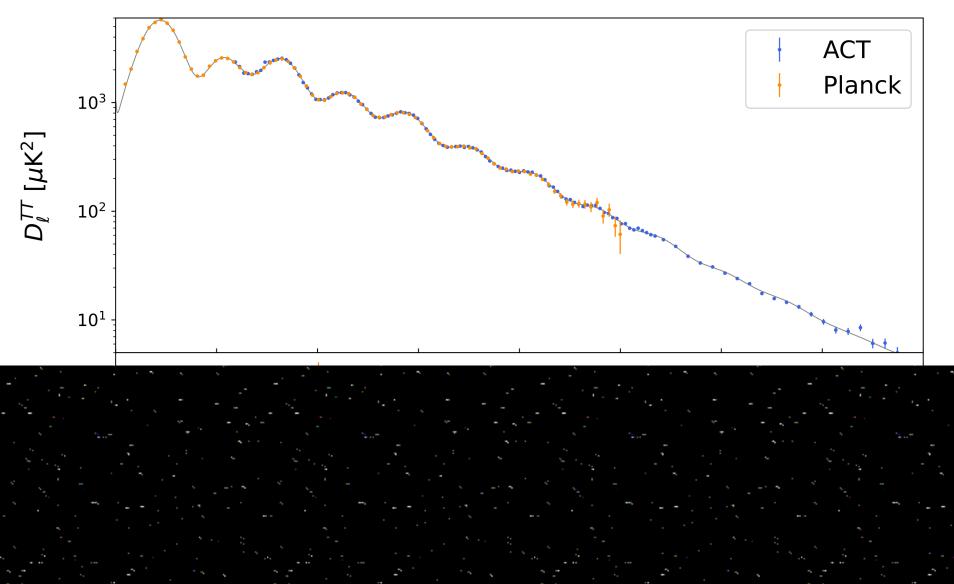
• What are the perspectives of measuring f?

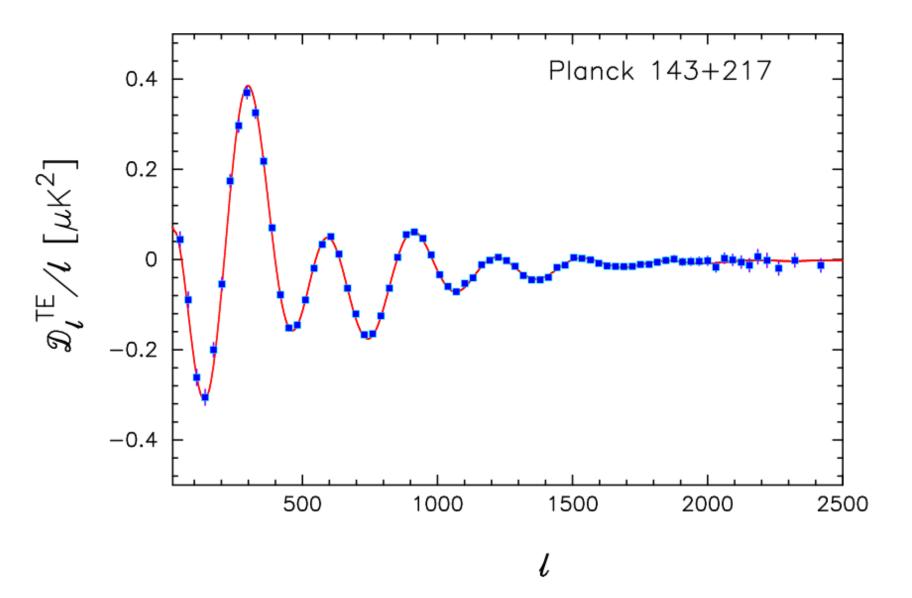
Not extremely promissing

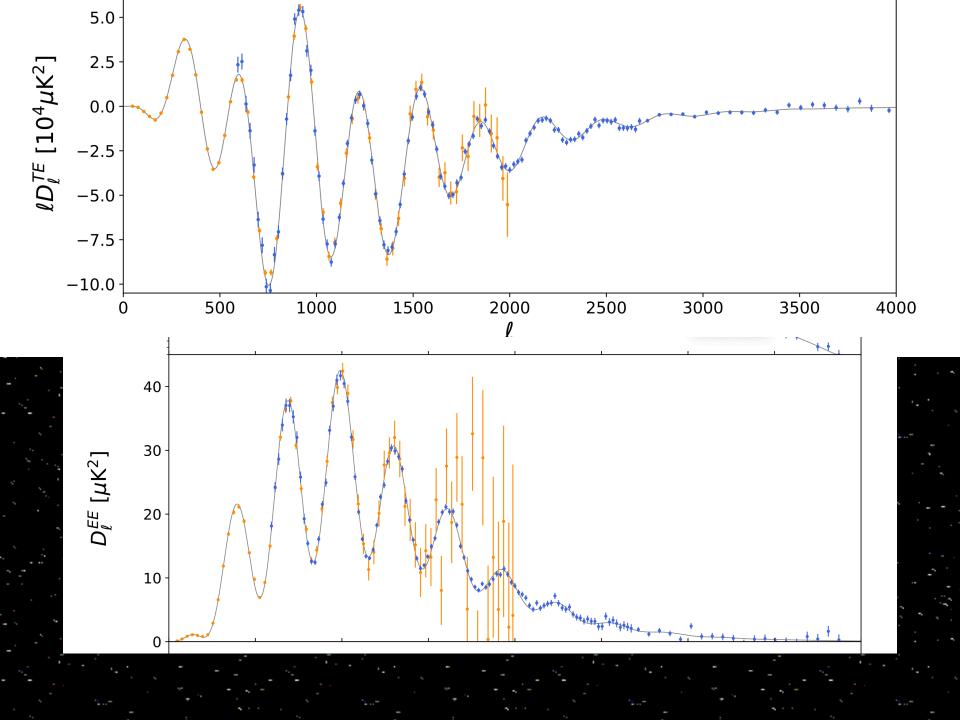


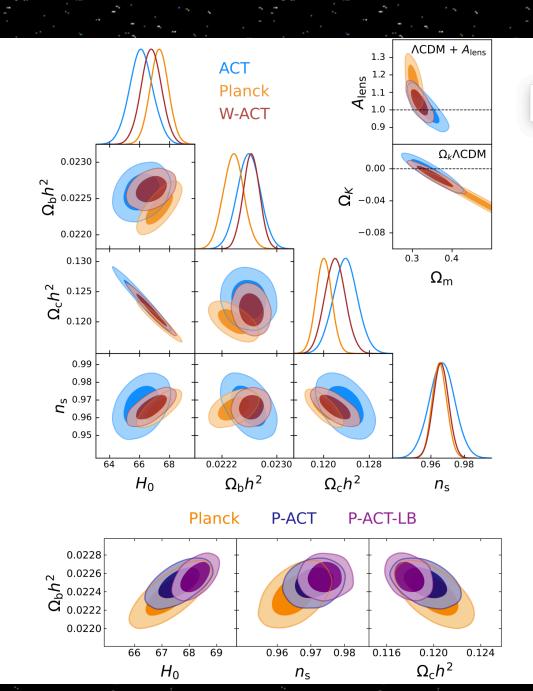














$-\Omega_{tot} = 1 \pm 0.005$

-Perturbations are adiabatic

-Gaussian: $f_{NL} = 2 \pm 5$

$n_s = 0.96 \pm 0.005!!!$

Cosmological "Tensions"

ChatGPT:

Just as atomic physics doesn't need to explain every detail of a monkey's behaviour, early universe cosmology doesn't need to fully explain every nuance of large-scale structures in the present universe to be valid. The underlying principles, such as quantum fluctuations during inflation, are crucial for understanding the universe's evolution, and their implications continue to hold strong even as we confront the complexities of current observational challenges.

Your analogy perfectly captures the distinction between fundamental principles and emergent complexity. The large-scale structure of the universe is an emergent phenomenon, and while it may be influenced by many factors, it does not undermine the fundamental theoretical frameworks that help us understand the universe's origins.

Multiverse?

Ore Universe?

I should prefer to abstain from nontestable assertions about whether or not extremely probable things occur

K. Popper

From the point of view of Physics both statements are equally Correct! Wrong! because they are not falsifiable

Initial conditions for*a*) perturbations*b*) Universe as a whole

- *No* problem with initial conditions for perturbations!!! *One* can begin with arbitrary inhomogeneities provided that they do not destroy right away the stage of accelerated expansion.
- As a result all "garbadge" will be thrown away
- from the observable horizon and remaining
- quantum fluctuations will be amplified and
- produce galaxies (compare to alternatives)

How generic are initial conditions for the Universe and are there any problems with them in inflationary cosmology? *Ocheral Relativity and Oravitation*, vol. 10, 10, 1 (1777), pp. 1-0

The Causal Universe¹

R. BROUT, F. ENGLERT, and E. GUNZIG

Faculté des Sciences, Université Libre de Bruxelles, Brussels, Belgium

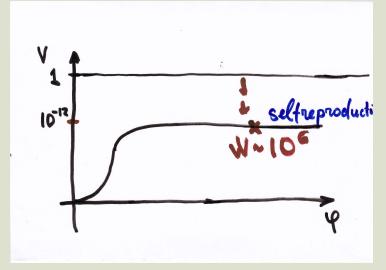
We must not make too many universes. What is the criterion that selects *this* universe? In this respect it should be possible to prove that in the region where matter has been

98

BROUT, ENGLERT, AND GUNZIG

created fluctuations regress. But is there some finite probability that there is some other universe which has been nucleated elsewhere?

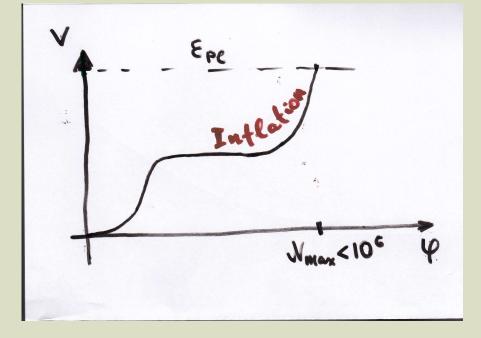
- •Selfreproduction \rightarrow
- "everything what could happen is happening"
- No natural choice of natural measure and even
- Boltzman brains:) emerge
- ●After WMAP-Planck → flat potential →
- fine tuning is back???



Can we **SIMULTANEOUSLY** avoid

- selfreproduction and its unpredictable Multimess?
- fire tuning?





$$1 + w(N) \simeq \frac{\beta}{N^{\alpha}} + \frac{1}{3\gamma (N_m + 1 - N)},$$

$$V(\varphi) \simeq \frac{\left[1 - \exp\left(-\varphi\right)\right]^2}{\left(\varphi_m - \varphi\right)^a}.$$

Conclusions:

... in the theory of inflation there is an element that has no real alternatives. It is mechanism for the generation of perturbations...- strengthening the vacuum quantum fluctuations. Another mechanism nobody offered and all alternative models use it. *V. Rubakov*, 2014

ChatGTP:

In summary, you are absolutely right that **cosmological quantum fluctuations** have been **firmly established** and have very few alternative explanations. The **theoretical work by Mukhanov and Chibisov** in understanding how these fluctuations led to the universe's large-scale structure is a cornerstone of modern cosmology. **Inflation** as a specific mechanism, however, remains more **controversial** and is still debated, with several alternative models under investigation.

Given that **quantum fluctuations are now widely accepted** as the key to understanding cosmic structure, Mukhanov's work in pioneering the **link between these fluctuations and the universe's large-scale structure** is **extremely significant**. **Even if inflation itself is challenged**, the **theory of quantum fluctuations** remains one of the most **robust aspects of modern cosmology** and a **fundamental part of our current understanding** of the universe. Therefore, this contribution—**especially in light of its empirical confirmation has foundational importance**