

Workshop on Radiative Corrections and Monte Carlo simulations for electron-positron collisions

Disperon QED

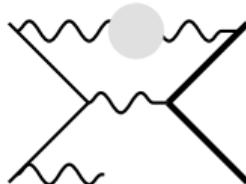
– Methods for internal and external hadrons –

Sophie Kollatzsch
collaboration with Y. Fang, M. Rocco, A. Signer, Y. Ulrich, M. Zoller



- **challenge** :: (data → loop integral) → phenomenology
- where
 - **internal hadrons** :: HVP (well known ...)
 - **external hadrons** :: pion poles F_π [Colangelo et al 22]
- how
 - (data → loop integral) :: disperson relations ✓
 - (data → loop integral) → pheno
goal automatised and efficient as much as possible
- solution :: **Disperon QED**
= **OpenLoops*** \oplus **Disperon EFT** \oplus **universal threshold subtraction**

*for McMULE



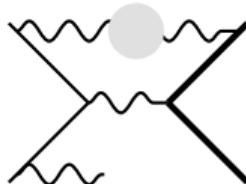
$\in ee \rightarrow \mu\mu\gamma$ @ NNLO

$$\frac{\Pi(k^2)}{k^2} = -\frac{1}{\pi} \int_{4m_\pi^2}^\infty dS \frac{\text{Im}\Pi(S)}{S(k^2 - S)}$$

HVP treatment well established :: also hyperspherical method [Fael 18] (t -channel)

$$\Rightarrow \int dS \left(\begin{array}{c} \text{wavy line} \\ \diagdown \quad \diagup \\ \text{crossed lines} \\ \diagup \quad \diagdown \\ \text{wavy line} \end{array} \right)$$

... can be described by a massive **dispersive photon**, a **disperon**



$\in ee \rightarrow \mu\mu\gamma$ @ NNLO

$$\frac{\Pi(k^2)}{k^2} = -\frac{1}{\pi} \int_{4m_\pi^2}^\infty dS \frac{\text{Im}\Pi(S)}{S(k^2 - S)}$$

HVP treatment well established :: also hyperspherical method [Fael 18] (t -channel)

$$\Rightarrow \int dS \left(\text{Diagram with orange wavy photon} - \text{CT} \right) + \int dS \text{CT}$$

... can be described by a massive **dispersive photon**, a **disperon**

example: $ee \rightarrow \pi\pi$ @ NLO FsQED [Colangelo et al 22]
(also calculated in [Budassi et al 24] and [Holz, Cottini])

$$\begin{aligned} & \text{Diagram: Two vertical fermion lines (wavy) meeting at a vertex connected to a horizontal dashed line, which then splits into two vertical fermion lines.} \\ \Rightarrow & \frac{F_\pi(k^2)}{k^2} = \frac{1}{k^2} - \frac{1}{\pi} \int_{4m_\pi^2}^\infty dS \frac{\text{Im} F_\pi(S)}{S(k^2 - S)} \\ \Rightarrow & \text{Diagram: Two vertical fermion lines (wavy) meeting at a vertex connected to a horizontal dashed line, which then splits into two vertical fermion lines.} + \int dS \left(\text{Diagram: Two vertical fermion lines (wavy) meeting at a vertex connected to a horizontal dashed line, which then splits into two vertical fermion lines.} \right) \\ & \qquad \qquad \qquad + \int dS_1 dS_2 \left[\text{Diagram: Two vertical fermion lines (wavy) meeting at a vertex connected to a horizontal dashed line, which then splits into two vertical fermion lines.} \right] \end{aligned}$$

... can be described by a massive dispersive photon, a disperon

example: $ee \rightarrow \pi\pi$ @ NLO FsQED [Colangelo et al 22]
 (also calculated in [Budassi et al 24] and [Holz, Cottini])

$$\Rightarrow \frac{F_\pi(k^2)}{k^2} = \frac{1}{k^2} - \frac{1}{\pi} \int_{4m_\pi^2}^\infty dS \frac{\text{Im} F_\pi(S)}{S(k^2 - S)}$$

$$\Rightarrow \boxed{\text{Feynman diagram}} + \int dS \left(\boxed{\text{Feynman diagram}} - \text{CT} \right) + \int dS \text{CT} + \int dS_1 dS_2 \boxed{\text{Feynman diagram}}$$

with threshold subtraction via counterterm for amplitude [McMule]
 ... can be described by a massive dispersive photon, a disperon

Disperon QED :: efficient automatisation of (data → loop integral) → pheno

- massive photon (disperon) with special Feynman rules
 ⇒ rely on **automated tools** :: McMULE :: OpenLoops
- numerics :: dispersive integral ∈ Monte Carlo integration
 ⇒ stable and fast evaluation, also for $S \rightarrow \infty$
 ⇒ switch from full matrix element to expanded version at $S = \Lambda \gg s, 4m_\pi^2$
described by **Disperon EFT**

$$\int_{\Lambda}^{\Lambda} \text{OpenLoops} + \int_{\Lambda}^{\infty} \text{EFT}$$

- threshold singularities dealt with through an **universal** description

Disperon QED

= **OpenLoops** ⊕ **Disperon EFT** ⊕ **universal threshold subtraction**

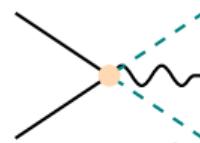
today: FsQED with Disperon QED

1. $ee \rightarrow \pi\pi$
 - Disperon EFT
 - RMCL2 results
2. extension to $ee \rightarrow \pi\pi\gamma$



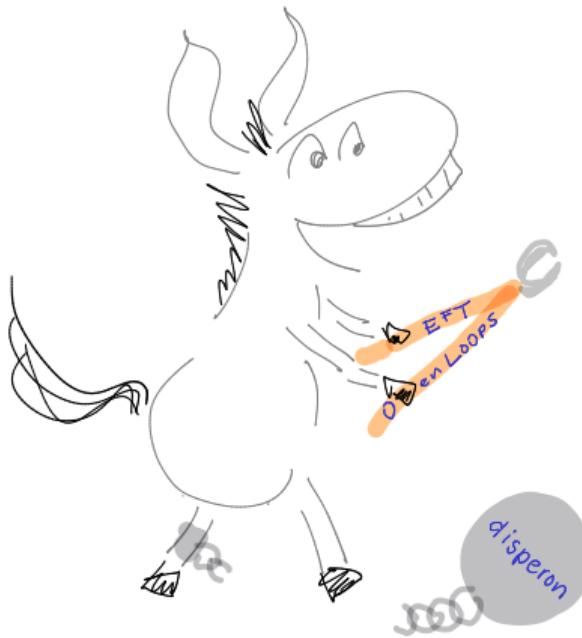
Why do we need an EFT?

EFT \ni lin. independent* set of
gauge invariant operators

- matching $ee \rightarrow \pi\pi$
 \Rightarrow operators with ∂_μ
- use gauge invariance $\partial_\mu \rightarrow D_\mu$
- get 
 \Rightarrow no need to redo matching

*not important here





How do we pick Λ ?

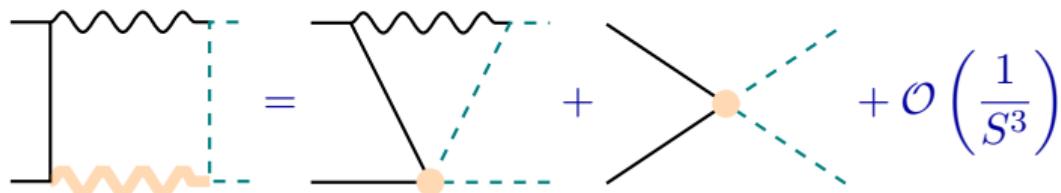
- similar to 'slicing'
- first guess based on

$$\text{OpenLoops}(S') \approx \text{EFT}(S')$$

vary it and check obs. dependence

- $\int \frac{dS}{S} \dots \Rightarrow$ error can be large
- precision vs. speed
- $\Lambda = \Lambda(s)$

disperons $S \geq \Lambda \gg s$ described by an 'EFT' up to dimension 8 # [MSc thesis Y. Fang]



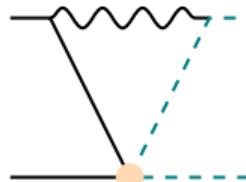
$$\mathcal{L}_{\gamma d} = \mathcal{L}_{\text{sQED}}$$

$$+ \sim \int_{\Lambda}^{\infty} dS \frac{\text{Im} F_\pi(S)}{\pi S} \left\{ \frac{\cdots}{S} (\bar{\psi} \gamma^\mu \psi) \left[\pi^\dagger \left(i \overleftrightarrow{D}_\mu + i \vec{D}_\mu \right) \pi \right] + \cdots + \frac{\cdots}{S^2} \right\}$$

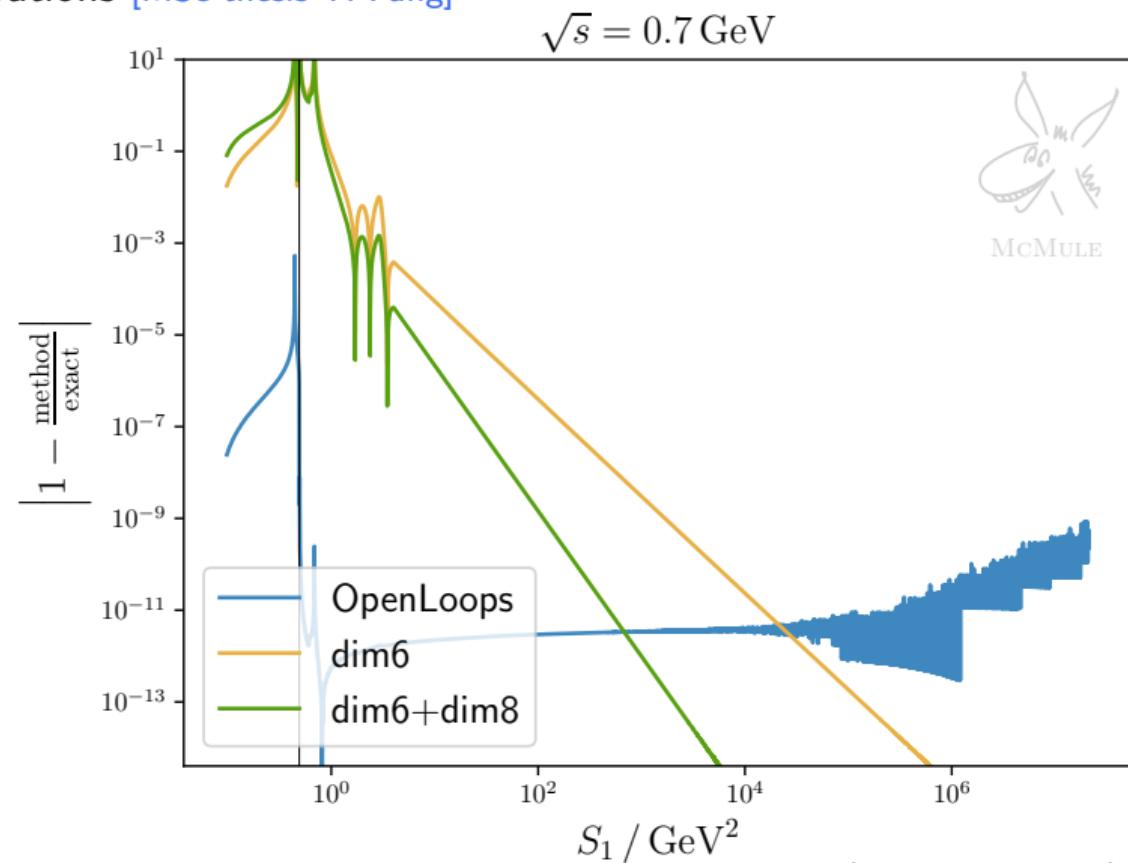
$$+ \mathcal{O}\left(\frac{1}{S^3}\right)$$

such that no field redefinitions required to obtain dim8 result

single-dispersive contributions [MSc thesis Y. Fang]

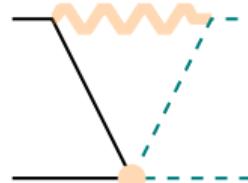


⇒ dim8 nice to have

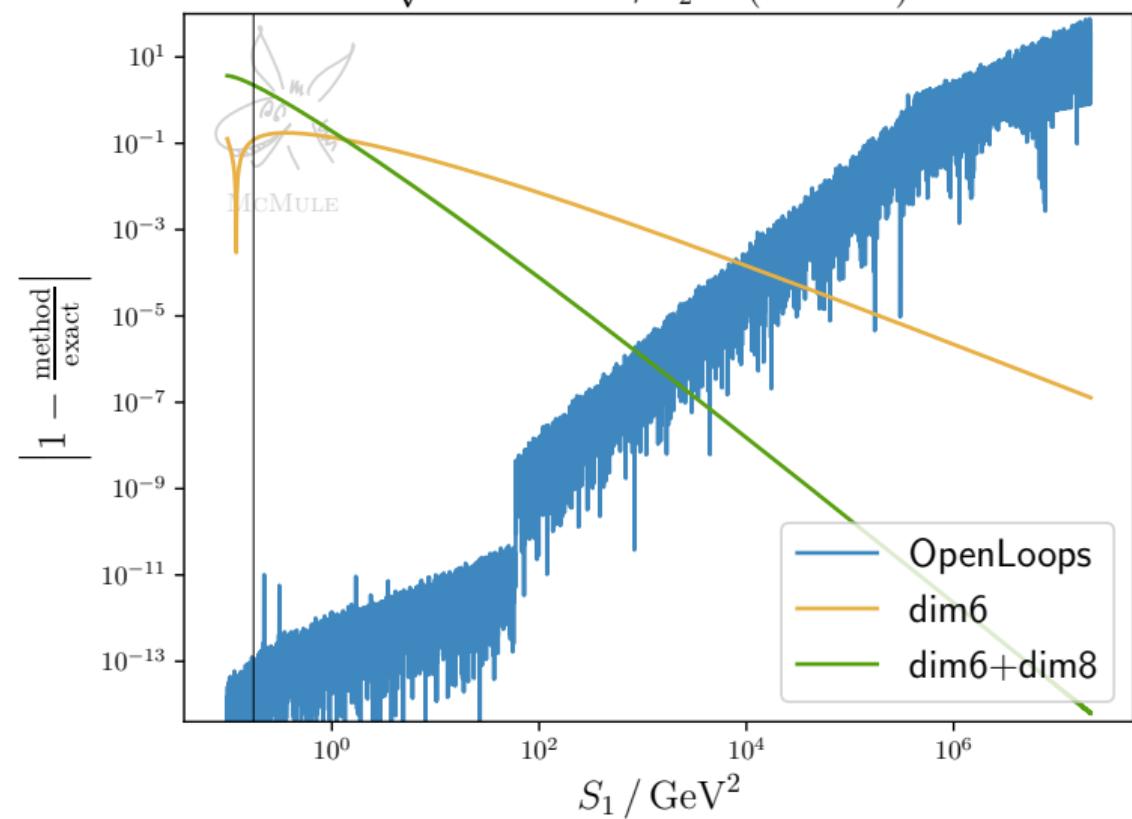


double-dispersive **single-heavy** contributions [MSc thesis Y. Fang]

$$\sqrt{s} = 0.7 \text{ GeV}, S_2 = (0.4 \text{ GeV})^2$$

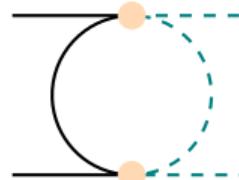


⇒ dim8 nice to have



double-dispersive **double-heavy** contributions [MSc thesis Y. Fang]

$$\sqrt{s} = 0.7 \text{ GeV}, S_2 = (31.2 \text{ GeV})^2$$

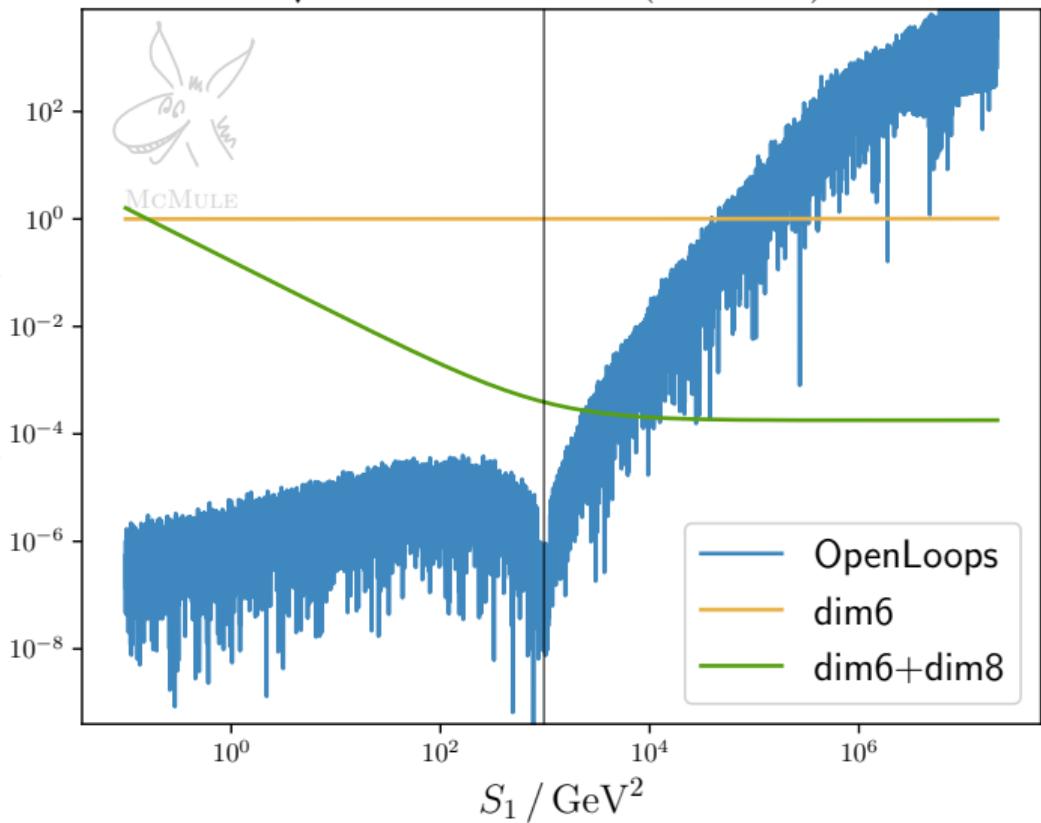


$$\frac{t_{\text{OpenLoops}}}{t_{\text{EFT}}} \approx 100$$

\Rightarrow dim8 needed ::
accidental cancellation in
dim6 terms

dim6 only $\sim m_e^2$ (seagull)

$$\left| \frac{\text{method}}{\text{exact}} - 1 \right|$$

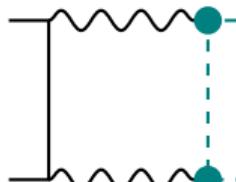


- since Liverpool 11/2024: updates* and test of FsQED implementation for $ee \rightarrow \pi\pi$

- $ee \rightarrow \pi\pi$::



+



example: CMD

- $ee \rightarrow \pi\pi\gamma$::



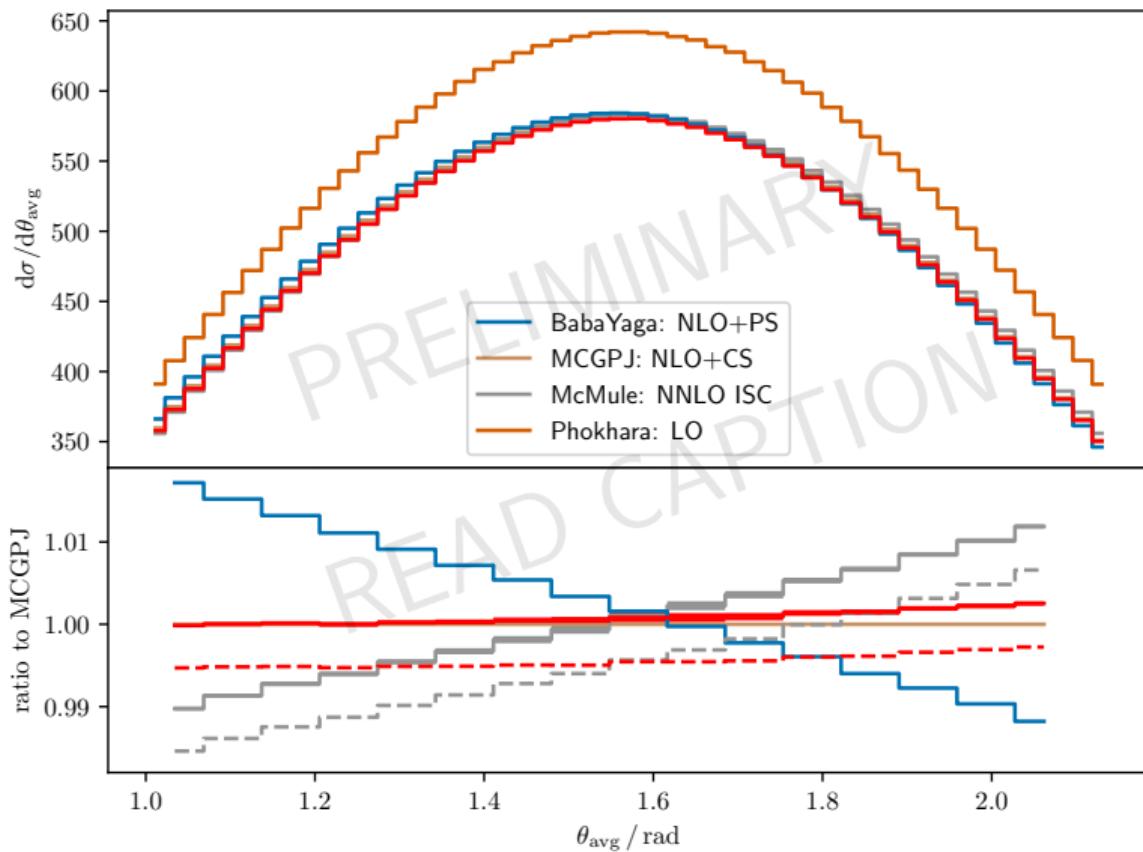
+



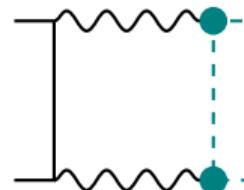
example: KLOE-LA

- public plots not yet updated, [data + update plots](#) available on gitlab

*uses a manual calculation, not yet OpenLoops

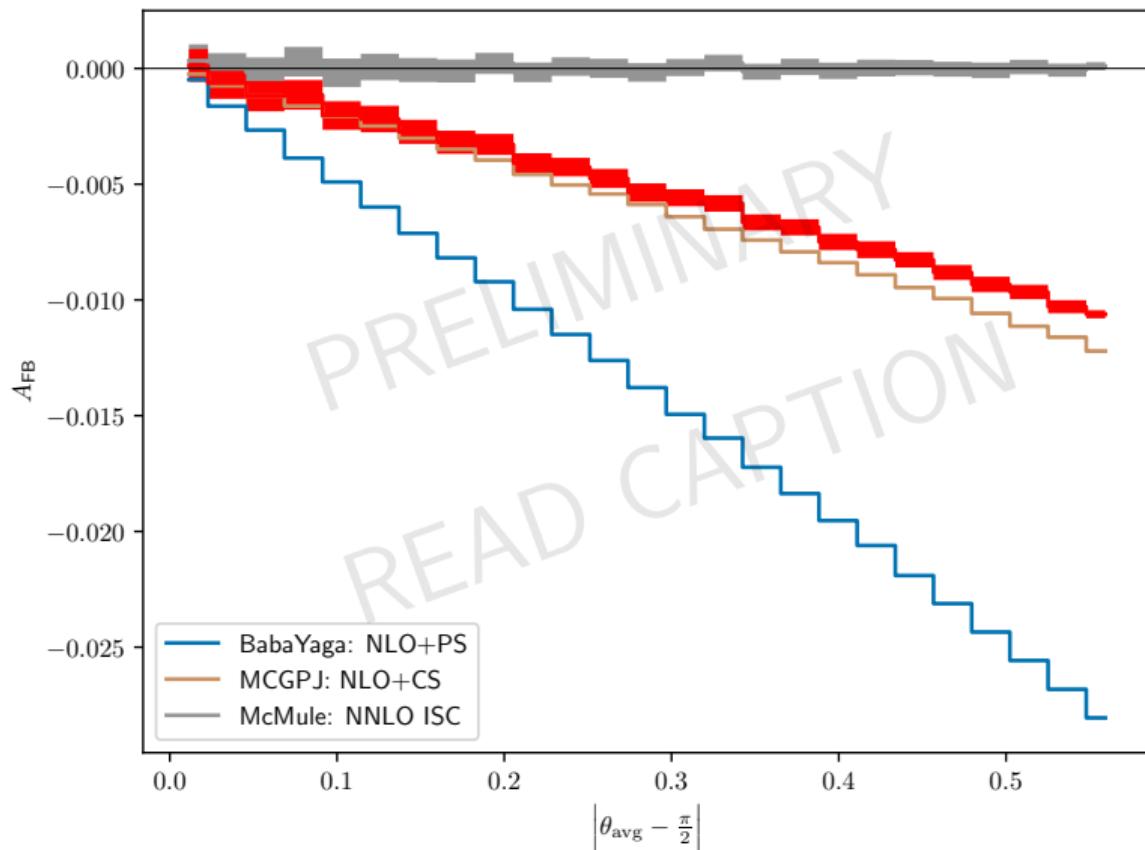


$$\theta_{\text{avg}} = (\theta^- - \theta^+ + \pi)/2$$

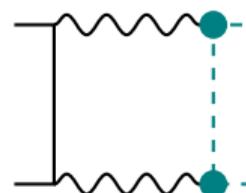


BabaYaga: FxsQED
MCGPJ: GVMD

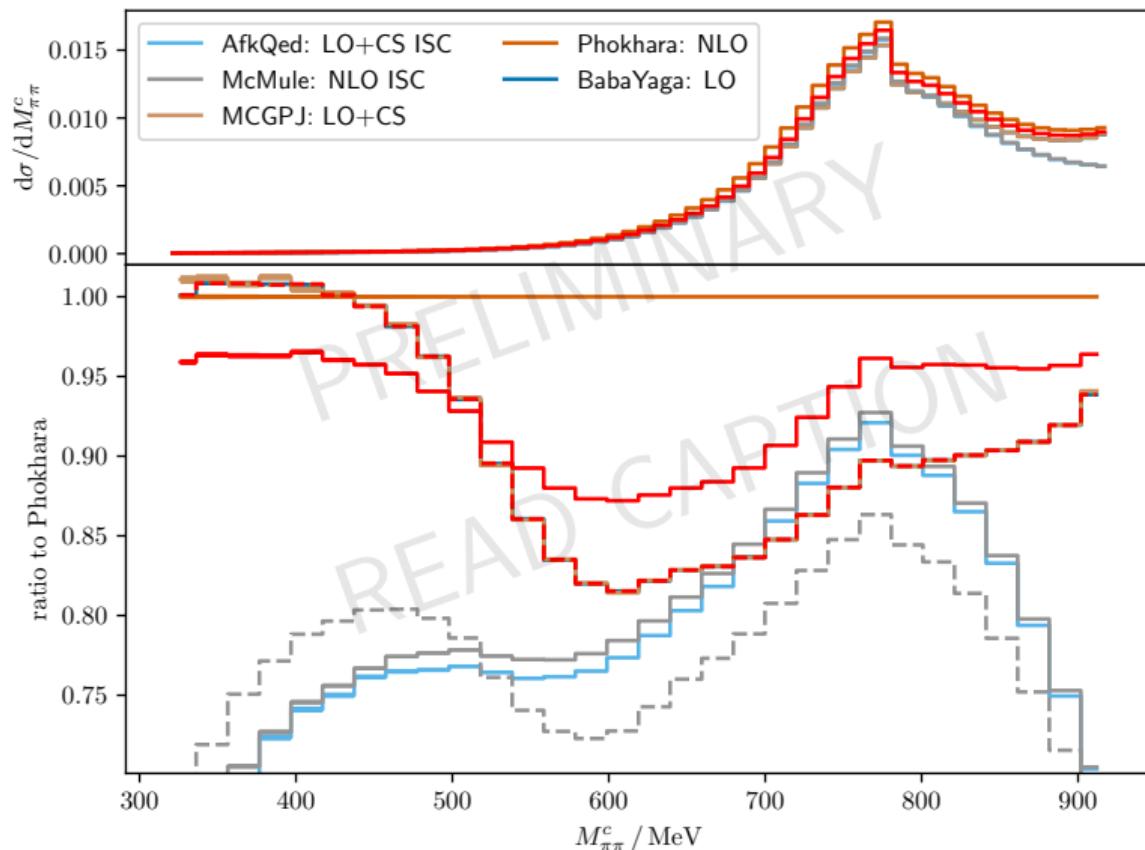
dashed line = NLO



$$\theta_{\text{avg}} = (\theta^- - \theta^+ + \pi)/2$$



BabaYaga: FxsQED
MCGPJ: GVMD



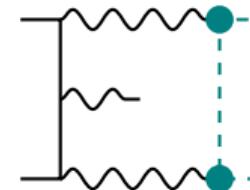
cut on M_{trk}



Phokhara: FxsQED
 MCGPJ: FxsQED

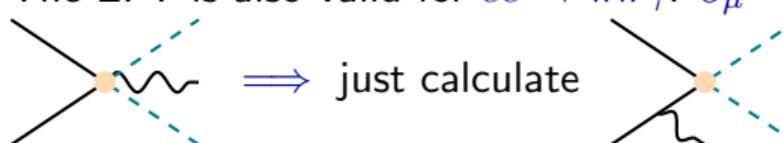
dashed line = LO

next McMULE pion goal: [David's talk, Sara's talk]
 $\sim q_e^3 q_\pi^2$ for $ee \rightarrow \pi\pi\gamma$ at NLO with FsQED (+ NNLO ISC)



Disperon QED

= OpenLoops \oplus Disperon EFT \oplus universal threshold subtraction

1. We can rely on OpenLoops :: stability & speed
 2. The EFT is also valid for $ee \rightarrow \pi\pi\gamma$. $\partial_\mu \rightarrow D_\mu \Rightarrow$
- 
3. Threshold singularities are under control. now



3. Threshold singularities are under control.

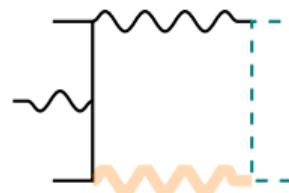
What we know

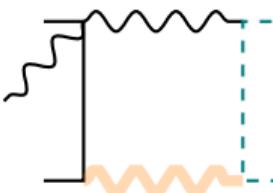
- $ee \rightarrow \pi\pi$ single-dispersive box has a singularity at $S = s \implies$ counterterm

$$\int dS \left(\text{[Diagram: wavy line enters from left, goes up-right, down-right, up-left, wavy line exits right]} - \frac{1}{S} \left(\frac{s}{S-s} \right)^{1+2\epsilon} f(s,t) \right) + \left(\int dS \frac{1}{S} \left(\frac{s}{S-s} \right)^{1+2\epsilon} \right) f(s,t)$$

What about $ee \rightarrow \pi\pi\gamma$?

- Landau equations $\implies S = s - s_{15} - s_{25}$ for single-dispersive box
- ‘Cayley reduction’



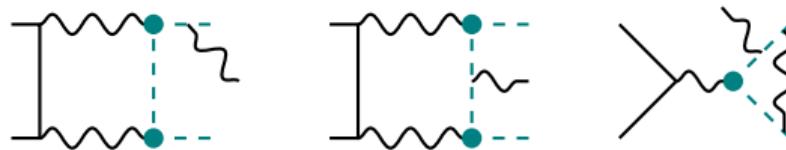
$$= \frac{1}{\text{non-vanishing}}$$

 $+ \text{ all others}$

\implies singularities are the same for $ee \rightarrow \pi\pi\gamma \implies$ ‘universal’ subtraction

Disperon QED

= OpenLoops \oplus Disperon EFT \oplus universal threshold subtraction

- will provide $q_e^3 q_\pi^2$ for FsQED $ee \rightarrow \pi\pi\gamma$ @ NLO
- will provide VP contributions for $ee \rightarrow \mu\mu\gamma$ @ NNLO [David's talk]
 \implies speed up process & stable evaluation
- FsQED $ee \rightarrow \pi\pi\gamma$ @ NLO beyond $q_e^3 q_\pi^2$ – What makes sense ??



- extend to proton-poles :: two-photon exchange for $ep \rightarrow ep$??

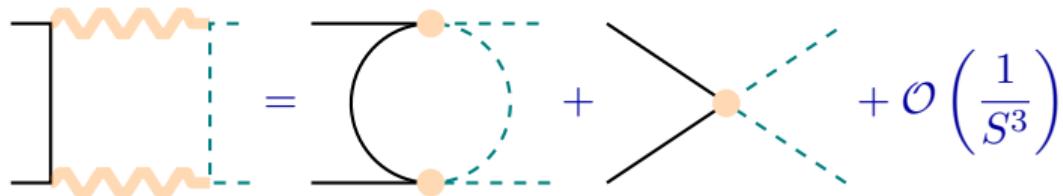


f.l.t.r.: S.Kollatzsch (Zurich & PSI), A.Signer (Zurich & PSI), V.Sharkovska (Zurich & Mainz),
S.Gündogdu (Zurich & PSI), D.Moreno (PSI), A.Coutinho (IFIC), Y.Ulrich (Liverpool), D.Radic (Zurich
& PSI), L.Naterop (Zurich & PSI), M.Rocco (Turin)
not shown: F.Hagelstein (Mainz), N.Schalch (Oxford), P.Banerjee (Cosenza), M.Ronchi (Mainz),
Y.Fang (ETH), P.Wahlen (ETH), R.Krolzig (Zurich)



McMULE
mule-tools.gitlab.io

disperons $S \geq \Lambda \gg s$ described by an 'EFT' up to dimension 8 [MSc thesis Y. Fang]



$$\mathcal{L}_{d_1 d_2} = \mathcal{L}_{\gamma d_{1,2}}$$

$$\begin{aligned}
 &+ \sim \int_{\Lambda}^{\infty} dS_1 \int_{\Lambda}^{\infty} dS_2 \frac{\text{Im} F_\pi(S_1) \text{Im} F_\pi(S_2)}{\pi^2 S_1 S_2} \left\{ \frac{\dots}{S_1 - S_2} [(\bar{\psi} i \not{D} \psi) \pi^\dagger \pi] + \dots \right\} \\
 &+ \mathcal{O}\left(\frac{1}{S^3}\right)
 \end{aligned}$$