Low Energy for KKMC and associated projects -Systematic Ambiguities what is to do

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• (1) KKMC for $e^+e^- \rightarrow l\bar{l}n\gamma$, $l = \mu, \tau, \nu, q$ (for τ 's with decays) was developed over several decades, starting from low energies, but since long focus mostly on high energy applications.

• (2) Program is useful for low energy applications. Some aspects, technical and theoretical are not developed to the point, mostly because of man power issues.

- (3) Let me present basic aspects first ...
- (4) ...and new applications later.
- (5) Long term man power perspective looks somehow better these days.
- (6) Talk plan: (i) Basic construction elements. (ii) Tests and low energies limitations. (iii) Versions and distributions (iv) New applications.

Funded in part by Narodowe Centrum Nauki, Poland, grant No. 2023/50/A/ST2/00224

• Separation into τ production and decay and spin algorithm.

• QED exclusive exponentiation. Amplitudes are constructed from parts. Eikonal part, correction obtained from first order amplitudes, correction obtained from second order amplitude ... etc. That required lots of work.

• Separating out: bremsstrahlung part of QED. Separating from the rest short scale interaction part: electroweak effects, at low energies that means $\alpha_{QED}(s)$ valid at first order, analytic and anti analytic properties of field theory results.

- Separating out New Physics effects, including interfering ones.
- Phase space parametrization, crude distribution, all possible because of conformal symmetry.
- Hopefully Alan presented some of these points.



• Because narrow τ width approximation can be obviously used for phase space, cross-section for the process $f\bar{f} \to \tau^+ \tau^- Y$; $\tau^+ \to X^+ \bar{\nu}$; $\tau^- \to X^- \nu$ reads:

$$d\sigma = \sum_{spin} |\mathcal{M}|^2 d\Omega = \sum_{spin} |\mathcal{M}|^2 d\Omega_{prod} \ d\Omega_{\tau^+} \ d\Omega_{\tau^-}$$

- This formalism is fine, but, e.g. for 20 τ decay channels we would have 400 distinct processes. Also picture of production and decay are mixed.
- Below only τ spin indices are explicitly written:

$$\mathcal{M} = \sum_{\lambda_1\lambda_2=1}^2 \mathcal{M}_{\lambda_1\lambda_2}^{prod} \; \mathcal{M}_{\lambda_1}^{ au^+} \mathcal{M}_{\lambda_2}^{ au^-}$$

• Cross section can be rewritten into core formula of spin algorithms

$$d\sigma = \left(\sum_{spin} |\mathcal{M}^{prod}|^2\right) \left(\sum_{spin} |\mathcal{M}^{\tau^+}|^2\right) \left(\sum_{spin} |\mathcal{M}^{\tau^-}|^2\right) wt \ d\Omega_{prod} \ d\Omega_{\tau^+} \ d\Omega_{\tau^-}$$

• where

$$wt = \left(\sum_{i,j=0,3} R_{ij}h^i h^j\right)$$

$$R_{00} = 1, \quad \langle wt \rangle = 1, \quad 0 \le wt \le 4.$$

 R_{ij} can be calculated from $\mathcal{M}_{\lambda_1\lambda_2}$ by contraction with Pauli σ^i matrices and similarly h^i , h^j respectively from \mathcal{M}^{τ^+} and \mathcal{M}^{τ^-} .

• Bell inequalities tell us that it is impossible to rewrite wt in the following form

$$wt \neq \left(\sum_{i,j=0,3} R_i^A h^i\right) \left(\sum_{i,j=0,3} R_j^B h^j\right)$$

that means it is impossible to generate first τ^+ and τ^- first in some given 'quantum state' and later perform separately decays of τ^+ and τ^-

- Approx. only! NO each τ spin in HepMC. Fatal: anom. dipole mom. pheno.

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- Phase-space Monte Carlo module producing "raw events".
- Library of models for provides input for "model weight"
- Useful for any application, not only auproduction/decay.
- Ratios of matrix elements squared define weight (probability) that event could be of model B if generated with mode A.
- Convenient for Machine Learning too.
- No compromises on precision are required.

Ref. frames for spin; production, decay. Geometry of QED amplitudes optimized



Figure 2

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- General idea: to identify in amplitudes, with the help of gauge invariance structures responsible later for phase-space enhancements: collinear-soft etc. This is fundamental, specially from the point of view of Monte Carlo algorithm construction.
- Discussions with **Shimizu-san** were important.
- Z. Was Gauge invariance, infrared / collinear singularities and tree level matrix element for e+ e- —> nu(e) anti-nu(e) gamma gamma Eur.Phys.J. C44 (2005) 489,
- A. van Hameren, Z. Was, *Gauge invariant sub-structures of tree-level double-emission exact QCD spin amplitudes*, Eur.Phys.J. C61 (2009) 33
- Also in this case algebraic manipulation methods were providing the reference calculations, necessary to cross check results.
- I was not able to find patterns automatically, but algebraic progams were essential for checks.
- Only some of the patterns appear naturally. Feynman diagrams 1 and 2 combined (next slide) are the complete amplitude for $\nu_{\mu}\bar{\nu}_{\mu}$ production.

Figure 1: The Feynman diagrams for $e^+e^- \rightarrow \bar{\nu}_e \nu_e \gamma$.





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- The first two diagrams represent initial state QED bremsstrahlun amplitudes for $\nu_{\mu}\bar{\nu}_{\mu}$ pair production. It can be divided into parts, corresponding to β_0, β_1 of Yennie-Frautshi-Suura exponentiation.
- Can separation be expanded to other cases, to higher orders, to terms of different singularities/enhancements?
- The answer seem to be yes, provided we stay in QED regime
- But it comes only after lots of effort.
- Many of NN-calculations do not respect (or can not respect) such separation into parts.
- It is most important for real emissions: algorithms design. Virtual corrections: broad aspect of separation into parts.
- The case is **testing limits** of QED YFS reorganization of perturbation expansion.
- Holds at second order too, but then internal charged scalar lines contribute. Is it QED?
- Virtual mixed virtual corrections were challenging too. But expansion around contact interaction for t-channel W's worked fine.
- Important: red part of the second page equation clearly identifiable. Essential feature used for algorithm design.

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$$\begin{split} \mathcal{M}_{1\{I\}} \begin{pmatrix} p_{k_{1}} \\ \lambda \sigma_{1} \end{pmatrix} &= \mathcal{M}^{0} + \mathcal{M}^{1} + \mathcal{M}^{2} + \mathcal{M}^{3} \\ \mathcal{M}^{0} &= eQ_{e} \ \bar{v}(p_{b},\lambda_{b}) \ \mathbf{M}_{\{I\}}^{bd} \ \frac{p_{a} + m - k_{1}}{-2k_{1}p_{a}} \ \epsilon_{\sigma_{1}}^{*}(k_{1}) \ u(p_{a},\lambda_{a}) \\ &+ eQ_{e} \ \bar{v}(p_{b},\lambda_{b}) \ \epsilon_{\sigma_{1}}^{*}(k_{1}) \ \frac{-p_{b} + m + k_{1}}{-2k_{1}p_{b}} \ \mathbf{M}_{\{I\}}^{ac} \ u(p_{a},\lambda_{a}) \\ \mathcal{M}^{1} &= \mathcal{M}^{1'} + \mathcal{M}^{1''} \\ \mathcal{M}^{1'} &= +e \ \bar{v}(p_{b},\lambda_{b}) \ \mathbf{M}_{\{I\}}^{bd,ac} \ u(p_{a},\lambda_{a})\epsilon_{\sigma_{1}}^{*}(k_{1}) \cdot (p_{c} - p_{a}) \frac{1}{t_{a} - M_{W}^{2}} \frac{1}{t_{b} - M_{W}^{2}}, \\ \mathcal{M}^{1''} &= +e \ \bar{v}(p_{b},\lambda_{b}) \ \mathbf{M}_{\{I\}}^{bd,ac} \ u(p_{a},\lambda_{a})\epsilon_{\sigma_{1}}^{*}(k_{1}) \cdot (p_{b} - p_{d}) \frac{1}{t_{a} - M_{W}^{2}} \frac{1}{t_{b} - M_{W}^{2}}, \\ \mathcal{M}^{2} &= +e \ \bar{v}(p_{b},\lambda_{b})g_{\lambda_{b},\lambda_{d}}^{We\nu} \ \epsilon_{\sigma_{1}}^{*}(k_{1}) \ v(p_{d},\lambda_{d})\bar{u}(p_{c},\lambda_{c})g_{\lambda_{c},\lambda_{a}}^{We\nu} \ k_{1} \ u(p_{a},\lambda_{a}) \frac{1}{t_{a} - M_{W}^{2}} \frac{1}{t_{b} - M_{W}^{2}}, \\ \mathcal{M}^{3} &= -e \ \bar{v}(p_{b},\lambda_{b})g_{\lambda_{b},\lambda_{d}}^{We\nu} \ k_{1} \ v(p_{d},\lambda_{d})\bar{u}(p_{c},\lambda_{c})g_{\lambda_{c},\lambda_{a}}^{We\nu} \ \epsilon_{\sigma_{1}}^{*}(k_{1}) \ u(p_{a},\lambda_{a}) \frac{1}{t_{a} - M_{W}^{2}} \frac{1}{t_{b} - M_{W}^{2}}, \\ (1) \end{array}$$

- Complete spin amplitude for e⁺e⁻ → ν
 _eν_eγ separated into six individually QED gauge invariant parts. Easy to check, replace photon polarization with its four-momentum. Each of the obtained parts has well defined physical interpretation. Eikonal universal infrared singular is in red. We touch essentials of pert. expansion reordering.
- It is also easy to verify that the gauge invariance of each part can be preserved to the case of the extrapolation, when because of additional photons, condition $p_a + p_b = p_c + p_d + k_1$ is not valid.

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- KKMC/Tauola/Photos all have exact complete phase space under control. That is good news for systematic.
- Ambiguities require studies of matrix elements.
- However, as off-shelf results are often partly intergated and not immediate to separate into parts they can not be used without costly adaptations...
- Another problem. For consistency, often irrelevant at low energies some approximatrions were used: same treatment of hard γ^* and Z imposed $\gamma \gamma$ box diagram approximations in KKMC.
- For photos MC it was phase space Jacobians matching of multibranch could not be better than order of ME applied. Expansion in curvature (Ricci tensor) must be consistent. Hard learned lesson.
- The path for ambiguities is clear but long and bumpy...

In previous slide I have show what is needed for some contributions.

Tests and low energies limitations

History:

- KORALB, KORALZ and predictions for many observable
- semi analytical and semi-analytical plus (structure functions calculations for selected observables, plus fixed order MC corrections).
- Low Energies (potential issues):
- Role of mass terms in radiative corrections. Important for radiative return configurations and box diagrams.
- radiative return logarithms and final state mass (phase space) effects.
- Pretabulation of electroweak corrections and what about $\alpha_{QED}(s)$, is it taken into account with sufficient granularity at low energies?
- Technical parameters may require different set-ups.
- Expertise of people, tests, man power, training.

Versions and their distributions

 \rightarrow Published versions of KKMC-f77 and KKMC-C++

 \rightarrow KKMC-F77 BASF2 of Belle collaboration validated $\alpha_{QED}(s)$ at 10 GeV CMS energies and below as well.

 \rightarrow https://github.com/KrakowHEPSoft/KKMCee kkmc@uj.edu.pl At present I am not involved in servicing this version. But I have it (after some difficulties) running. On the way I had encounter annoying problems due to some versions of root : old were ok, new were ok, intermediate were messing up event weights.

- Some tests are running fine, but many are waiting to be transferred from FORTRAN.
- Some known, and fixed issues, like spin entries in HepMC3 seem not to be transmitted to distribution.

- If anomaly of $(g-2)_{\mu}$ is a reality then there should be an impact elsewhere too.
- New interactions are often proportional to the lepton mass.
- Then, why not in τ leptons?
- Also, why not anomalous electric/magnetic dipole moments?
- How this could manifest itself in observables?
- We have implemented such extra interactions first for KKMC MC at Belle 2 energies Phys.Rev.D 106 (2022) 11, 113010
- Then for KKMC at FCC energies 2407.17282
- Recently (2407.17282 and next week code semi public) for TauSpinner and LHC applications.

From g-2 to τ lepton decays

In collaboration with by **A. Korchin**^{†,*}

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- (1) Anomalous dipole moment $(g-2)_{\mu}$ of interest \rightarrow theory-experiment tensions.
- (i) To resolve: (i) measure better, (ii) calculate better; revise its input,
 (iii) discovery. experiments need new theory, then theory other experiments, which...
- (3) My experience goes with (ii), in particular:

- (4) rad. return evts. for $\alpha_{QED}(s)$? predictions for $e^+e^- \rightarrow \pi^+\pi^-(l^+l^-)n\gamma$ need improvements. at work \rightarrow man power, motivation from (g-2)_µ helps.

- (5) Also for (g-2) ambiguities: help from $\tau \to \nu \ hadrons$ for $\alpha_{QED}(s)$ too (after isospin rotation and optical theorem).

- (6) If anomalous $(g-2)_{\mu}$, then electric magnetic dipole moments in τ 's? • (7) I can not ignore kitchen for τ production and decay, also how bremsstrahlung (i.e. QED) separates from hard interactions.

• Beyond Standard Model anomalous dipole moments may come from New Physics particles in loop corrections to $Z(\gamma) - \mu^+ - \mu^-$ vertex. If seen, that would bring fundamental change.

• Tension between theory predictions and experimental measurement for anomalous magnetic dipole moment $(g-2)_{\mu}$ is at the level of 5σ : Fermilab Muon g-2 Experiment, Phys Rev Lett **131** (2023) 161802. If it will survive time test \rightarrow change in our understanding of physics foundations.

• Measurements (high precision): $a_{\mu}(\text{FNAL}) = 116\,592\,055(24)\cdot10^{-11}(0.20 \text{ ppm}).$ New experimental world average: $a_{\mu}(\exp) = 116\,592\,059(22)\cdot10^{-11}$ (0.19 ppm).

• For the theory side: Muon g-2 Theory Initiative: T. Aoyama et al., *The anomalous magnetic moment of the muon in the Standard Model*, Phys. Rep. 887, 1 (2020).

QED multi loop calculations are impressive: T. Aoyama, M. Hayakawa, T. Kinoshita, and M. Nio, Complete Tenth-Order QED Contribution to the Muon g-2, Phys. Rev. Lett. 109, 111808 (2012). But this is not the end of the story...

AMBIGUITIES is my life

I will address only some of them. Those where I have something to say.

Systematic of KKMC $e^+e^- \to \mu^+\mu^-(or\ hadrons)\ n\gamma$ and τ decay Monte Carlos, slide no. 12.

In particular of $e^+e^- \rightarrow \mu^+\mu^- n\gamma$, with hard γ 's, and τ decay Monte Carlos spectral functions and what will change if final state is $\pi^+\pi^-$.

- Because of systematic ambiguities for g-2 interpretation, there is an interest to extend data pool for dispersion relations used in $\alpha_{QED}(s)$ evaluation. That means $e^+e^- \rightarrow hadrons$ total cross section as function of s.
- \bullet Belle data can be useful, the KKMC radiative return to small invariant masses of $\mu\text{-pair}$ can be used if...
- \rightarrow Muon mass has to be changed to pion mass.
- \rightarrow Its PDGid changed,
- ightarrow Born level cross section re-weighted from fermion to scalars
- ightarrow effect of intermediate resonances imprinted to the weight as well.
- How well it can work? Theoretical ambiguities will be larger than for μ .
- Anyway another helpful tool. Cross check with Phokara.
- Implementation prepared by Jadach is not easy to recover.

- Bremsstrahlung and its virtual counter-part happen at smaller energies than hard interaction e.g. of $\gamma^* \rightarrow \tau^+ \tau^-$.
- That sounds simple, but requires effort to separate. Without proper evaluations of Standard Model predictions and their ambiguities:
- no input for $\alpha_{QED}(s)$ from radiative return events

• If New Physics would not separate consequences would be less severe: that would be New Physics anyway, but its nature would be more difficult to interpret.

- To do that one has to keep under control issues of factorization:
- scale of QED bremsstrahlung (with its virtual corrections) differ from scale of hard process.
- That is useful for intuition, but if ambiguities are to be taken into account, serious work on properties of spin amplitudes are needed.
- Work on semi inclusive quantities, is not enough. Thanks to the efforts, we can advocate Mustraal frame to separate hard and bremsstrahlung (parton level) interactions parts of phase-space and amplitudes.
- I will cover this topic, at a side of discussion on: turning $e^+e^- \rightarrow \mu^+\mu^- n\gamma$ Monte Carlo into one of $e^+e^- \rightarrow \pi^+\pi^- n\gamma$. It is of importance in itself though.

• Promising path \rightarrow use Mustraal frame. Minimize bias from lack of bremsstrahlung amplitudes and virtual corrections for π 's. At the time of LEP 1 Monte Carlo (1982) work, amplitudes for $e^+e^- \rightarrow \mu^+\mu^-\gamma$ were carefully studied.

• It was found (no approximation): fully differential cross section could be written in form of factorization of photon emission factor and corresponding phase space term of photon energy and its $\theta\phi$ angles, times the factor of the same structure as Born.

- Sum of two expressions of slightly different frames orientation was necessary.
- Formed incoherent sum, excellent for MOnte Carlo implementation.
- There is a theorem by Rondald Kleiss that such separation is quite universal and got first correction from α^2 terms (without any enhancement).
- We have collected numerous tests for that in context of spherical harmonics used nowdays in LHC data interpretation.
- The checks were for QCD and up to second order (two parton emissions).
- No need to repeat that work for $\pi^+\pi^-$: properties originate from Lorentz group structure: its layers with respect to rotation sub-group.

Figure 2: Mustraal frame $\cos \theta$ distributions for muons. Only hard bremsstrahlung events taken.

Test level 0 : No deformations with respect to Born level shape.



Figure 3: Mustraal frame $\cos \theta$ distributions for pions. re-weighting from muon sample. Only hard bremsstrahlung events taken. Reweighting of BASF2-lib events with user code.

Test level 0 : No deformations with respect to Born level shape.



• But what about FSR and IFI interference, for $\pi^+\pi^-$ final states?

• Message from scalar QED is not very encouraging. amplitudes for $\gamma^* \rightarrow \pi^+ \pi^- \gamma$ are complicated (0802.2182) have their own enhancements, which **DO NOT** match QED. Even QCD results (0802.2182) look better.

• Hope: feed back (fits, form factors) from experiments like BESIII may settle the matter.

• Not as straighforward and theoretical considerations inescapable?

• GOOD: eikonal part of amplitdue separates out. Complete first order contribution to final state bremsstrahlung: $\gamma^* \to \pi^+ \pi^- \gamma$, can be unambiguously added through event reweighting.

- SURPRIZE: this weight is becoming large when $\pi^+\pi^-$ pair invariant mass is large. No problem for Photos, but possibly points to ambiguity and in particular in interference with ISR.
- **BAD:** Radiative return ambiguity? Interference with ISR may be larger than I would bet before Pisa.

Outlook

- The purpose of my talk was to push some ideas forward and what is needed for that.
- It was not disciplined talk. Sorry for that.
- My aim was to underline directions, which in my opinion, are worth to follow.
- I have sketched also what are the main aspects of the activities now.
- I hope it may be useful.
- I promise to contribute, I hope to complete training of new bunch of contributors too.
- I hope they will then find path to continue, stay in the field etc. etc.
- Thanks for listening.

ps. In Cracow it is even colder than here. Last week we had +28 degrees.