## LEVERHULME TRUST



# **Progress on Soft Photon Resummation Implementation in Phokhara**

Jérémy Paltrinieri







- Importance of hadron cross-sections measurements in g-2 determination
- Radiative corrections studies must be undertaken in order to clear up the tensions between experiments
- From radiative return experiments like KLOE, this means taking care of NNLO and all-order exponentiation effects: many challenges!
- In Phokhara, pursing NNLO accuracy on one side and approximate all order accuracy in the soft photon regime on the other side, for radiative return processes such as  $e^+e^- \rightarrow \mu^+\mu^-\gamma$

- photon emission events and the emission of two real hard photons.
- Case of muons: complete NLO [arXiv:1312.3610]
- Case of pions: FxsQED for the pion treatment [Szymon Tracz PhD Thesis]
- Disclaimer: not discussing many other hadronic final states:  $3\pi$ ,  $4\pi$ , NN, KK...

 Phokhara simulates electron-positron annihilation into hadrons (and muons) at NLO accuracy. This includes virtual and soft photon corrections to one







- Current version: Phokhara v10.0: https://looptreeduality.csic.es/phokhara/ (October 2020)
- Last update: two remaining gauge-invariant contributions: FSRNLO & TVP
- Strong2020 paper [2410.22882] and website







# **Roadmap for Phokhara**

 $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$  @ NLO sQED for pions

CEEX formalism implemented in Phokhara, with the goal of implementing resummed  $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ 

valid at  $\mathcal{O}(\alpha^1)_{CEEX}$  in the language of KKMC

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# **Roadmap for Phokhara**



Treatment of pions: see Pau's talk (insertion of GVMD model in Phokhara)

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# **Roadmap for Phokhara**

## $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ @ NNLO See William's talk



Treatment of pions: see Pau's talk (insertion of GVMD model in Phokhara)

 $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$  @ NLO sQED for pions



# Resummed CEEX results for $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ Valid at $\hat{\mathcal{O}}(\alpha^2)_{CEEX}$ ?

CEEX formalism implemented in Phokhara, with the goal of implementing resummed  $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ 

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# **Towards CEEX in Phokhara**

- Sheds another light on missing higher-order effects
- YFS formalism: resummation of soft photons effects
- 2-loop integrals are notoriously hard: approximations provide insight!



CEEX Private			⊙ Unwatch 1	• 양 Fork 0 • ☆
main 👻 🖁 H Branch 💿 0 Tags	Q Go to file	t Add file 👻	<> Code -	About
jpaltrin updated README		0dc50b2 · 1 hour ago	🕚 18 Commits	Implementation of Soft Phot Resummation in Fortran
ANALYSIS	multiple histograms		19 hours ago	🛱 Readme
NOTES	added eikonal notebook		last week	-∿- Activity
RESULTS	cleaner environment		last week	⊙ 1 watching
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README			∅ :≡	No packages published Publish your first package

# Differential Distributions for $e^+e^- \rightarrow \mu^+\mu^-$ within the CEEX Framework

This Fortran code computes differential distributions for the process  $e^+e^- \rightarrow \mu^+\mu^-$ , taking into account lepton masses. It supports calculations at Leading Order (LO) and includes optional soft photon resummation, called Coherent Exclusive Exponentiation (CEEX), which is WIP.

### Packages No packages published Publish your first package Languages Python 40.4% Fortran 3 Mathematica 23.0% She Makefile 0.6%

Based on your tech stack

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$$\sigma^{(r)} = \frac{1}{\text{flux}(s)} \sum_{n=0}^{\infty} \int d\text{Lips}_{n+2}(p_a + p_b; p_c, p_d, k_1, \dots, k_n) \ \rho_{\text{CEEX}}^{(r)}(p_a, p_b, p_c, p_d, k_1, \dots, k_n)$$

 $ho_{ ext{CEEX}}^{(r)}(p_a,p_b,p_c,p_d,k_1,k_2,\ldots,k_n)$ 

$$\sum_{i,j,l,m=0}^{3} \hat{\varepsilon}_{a}^{i} \hat{\varepsilon}_{b}^{j} \sigma_{\lambda_{a}\bar{\lambda}_{a}}^{i} \sigma_{\lambda_{b}\bar{\lambda}_{b}}^{j} \mathfrak{M}$$

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# **Progress on CEEX in Fortran**

$$_{n}) = \frac{1}{n!} e^{Y(\Omega; p_{a}, \dots, p_{d})} \bar{\Theta}(\Omega) \sum_{\sigma_{i} = \pm 1} \sum_{\lambda_{i}, \bar{\lambda}_{i} = \pm 1}$$

 $\mathfrak{N}_{n}^{(r)}\left(\begin{smallmatrix}p k_{1} k_{2} \\ \lambda \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \left[\mathfrak{M}_{n}^{(r)}\left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{M}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \end{smallmatrix}\right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{1} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)} \left(\begin{smallmatrix}p k_{1} k_{2} \\ \bar{\lambda} \sigma_{2} \right) \cdot \left[\mathfrak{m}_{n}^{(r)$ 



$$\sigma^{(r)} = \frac{1}{\text{flux}(s)} \sum_{n=0}^{\infty} \int d\text{Lips}_{n+2}(p_a + p_b; p_c, p_d, k_1, \dots, k_n) \rho^{(r)}_{\text{CEEX}}(p_a, p_b, p_c, p_d, k_1, \dots, k_n)$$

Multi-Photon Phase-Space Generation

 $\rho_{\text{CEEX}}^{(r)}(p_a, p_b, p_c, p_d, k_1, k_2, \dots, k_n)$ 

$$\sum_{i,j,l,m=0}^{3} \hat{\varepsilon}_{a}^{i} \hat{\varepsilon}_{b}^{j} \sigma_{\lambda_{a}\bar{\lambda}_{a}}^{i} \sigma_{\lambda_{b}\bar{\lambda}_{b}}^{j} \mathfrak{M}$$

Jérémy Paltrinieri

# **Progress on CEEX in Fortran**

$$(n_i) = \frac{1}{n!} e^{Y(\Omega; p_a, \dots, p_d)} \bar{\Theta}(\Omega) \sum_{\sigma_i = \pm 1} \sum_{\lambda_i, \bar{\lambda}_i = \pm 1}$$

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YFS form factor

Multi-Photon Phase-Space Generation

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$$(x_n) = rac{1}{n!} e^{Y(\Omega; p_a, \dots, p_d)} \bar{\Theta}(\Omega) \sum_{\sigma_i = \pm 1} \sum_{\lambda_i, \bar{\lambda}_i = \pm 1}$$

Beta functions (IR safe Matrix Elements)



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# **Progress on CEEX in Fortran**



## Analysis Framework

$$\sigma_n) = \frac{1}{n!} e^{Y(\Omega; p_a, \dots, p_d)} \bar{\Theta}(\Omega) \sum_{\sigma_i = \pm 1} \sum_{\lambda_i, \bar{\lambda}_i = \pm 1} \sum_{\lambda$$

Beta functions (IR safe Matrix Elements)





# Progress on CEEX in Fortran

Multi-Photon Phase-Space Generation

Module PS

Multi-photon Phase-Space event generator is necessary for resummed results. Can be cross-checked with NLO generator in Phokhara.

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Multi-Photon Phase-Space Generation

Module PS

Multi-photon Phase-Space event generator is necessary for resummed results. Can be cross-checked with NLO generator in Phokhara.

### Input:

- n: integer, number of photons **Outputs**:

- momenta (modified in-place)
- weight: phase space weight (modified in-place)

### Step 1: Initialise incoming momenta

- Generate initial energy and momentum
- Add positron and electron momenta
- Update weight

### **Step 2: If no photons**

- Reconstruct final state muon-antimuon pair momenta

### **Step 3: If photons are emitted**

Repeat until a valid event is generated:

- For each photon:

- Generate massless photon momentum - Reconstruct final state muon-antimuon pair momenta - Check kinematic validity:

- Update weight

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# **Progress on CEEX in Fortran**

![](_page_16_Picture_26.jpeg)

![](_page_17_Picture_0.jpeg)

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# **Progress on CEEX in Fortran**

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0.500000	0.00000	0.000000	0.499975	0
0.500000	0.00000	0.00000	-0.499975	0
0.500000	0.051492	0.043576	-0.484031	0
0.500000	-0.051492	-0.043576	0.484031	0
Total mome	ntum			
0.000000	0.000000	0.000000	0.000000	

### **0-photon event**

E	px	ру	pz	 m2
0.500000 0.500000 0.278691 0.223409 0.131952 0.020522 0.047267 0.027761 0.167504	0.000000 0.000000 -0.087642 -0.027450 -0.018749 0.001769 0.023795 -0.010782 0.127600	0.000000 0.000000 0.088248 -0.008661 0.024650 -0.001243 -0.004575 -0.000331 -0.088527 -0.000550	0.499975 -0.499975 -0.225912 0.194729 -0.128266 -0.020407 0.040584 -0.025580 0.062760	0.000025 0.000025 0.011164 0.011164 -0.000000 0.000000 0.000000 -0.000000 -0.000000
0.102894  Total mome 0.000000	-0.008541  ntum 0.0000000	0.000000	-0.000000	

### 6-photon event

![](_page_17_Picture_33.jpeg)

# Progress on CEEX in Fortran

![](_page_18_Figure_1.jpeg)

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# YFS Form Factor

# Module FormFactor

# Jérémy Paltrinieri

### Input:

- pl: momenta particle l
- p2: momenta particle 2

### **Outputs**:

- YFS formfactor B(p1,p2)

### Virtual IR Function

Here we present the expression for the virtual part of the YFS for any two charged massive particles.

$$2\alpha \mathcal{R}B(p_{1},p_{2}) = -Z_{i}Z_{j}\theta_{i}\theta_{j}\frac{\alpha}{\pi} \bigg[ \bigg(\frac{1}{\rho}\ln\frac{\mu(1+\rho)}{m_{1}m_{2}} - 1\bigg)\ln\frac{m_{\gamma}^{2}}{m_{1}m_{2}} + \frac{\mu\rho}{s}\ln\frac{\mu(1+\rho)}{m_{1}m_{2}} + \frac{m_{1}^{2} - m_{2}^{2}}{2s}\ln\frac{m_{1}}{m_{2}} - 1 \\ + \frac{1}{\rho}\bigg(\pi^{2} - \frac{1}{2}\ln\frac{\mu(1+\rho)}{m_{1}^{2}}\ln\frac{\mu(1+\rho)}{m_{2}^{2}} - \frac{1}{2}\ln^{2}\frac{m_{1}^{2} + \mu(1+\rho)}{m_{2}^{2} + \mu(1+\rho)} - \operatorname{Li}_{2}(\zeta_{1}) - \operatorname{Li}_{2}(\zeta_{2})\bigg)\bigg],$$
(A.1)

where,

$$\mu = p_1 p_2, \quad s = 2\mu + m_1^2 + m_2^2, \quad \rho = \sqrt{1 - \left(\frac{m_1 m_2}{\mu}\right)^2}, \quad \zeta_i = \frac{2\mu\rho}{m_i^2 + \mu(1+\rho)}.$$
 (A.2)

### **Real IR Function**

Here we present an expression for the IR function  $\tilde{B}$  which corresponds to the emission of a real photon  $k \in \Omega$  from a dipole consisting of two charged massive particles  $p_1$  and  $p_2$ .

$$2\alpha \tilde{B}(p_1, p_2) = -Z_i Z_j \theta_i \theta_j \frac{\alpha}{\pi} \left[ \left( \frac{1}{\rho} \ln \frac{\mu(1+\rho)}{m_1 m_2} - 1 \right) \ln \frac{\omega}{m_\gamma^2} + \frac{1}{2\beta_1} \ln \frac{1+\beta_1}{1-\beta_1} + \frac{1}{2\beta_2} \ln \frac{1+\beta_2}{1-\beta_2} + \mu G(p_1, p_2) \right],$$
(A.3)

logarithms and dilogarithms,

$$G(p_1, p_2) = \frac{1}{\sqrt{(Q^2 + M^2)(Q^2 + \delta^2)}} \left[ \ln \frac{\sqrt{\Delta^2 + Q^2} - \Delta}{\sqrt{\Delta^2 + Q^2} + \Delta} \left[ \chi_{23}^{14}(\eta_1) - \chi_{23}^{14}(\eta_0) \right] + Y(\eta_1) - Y(\eta_0) \right],$$
(A.4)

# **Progress on CEEX in Fortran**

### [2203.10948] **Sherpa YFS**

where  $\beta_i = \frac{|\vec{p}_i|}{E_i}$  and  $\mu, \rho$  are defined in Eq. (A.2) and  $\omega$  is the momentum cut-off specifying  $\Omega$  in the frame  $\tilde{B}$  is to be evaluated in.  $G(p_1, p_2)$  is a complicated function that can be expressed as a combination of

# YFS Form Factor

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# Jérémy Paltrinieri

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- p2: momenta particle 2

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(A.1)

where.

$$\mu = p_1 p_2, \quad s = 2\mu + m_1^2 + m_2^2, \quad \rho = \sqrt{1 - \left(\frac{m_1 m_2}{\mu}\right)^2}, \quad \zeta_i = \frac{2\mu\rho}{m_i^2 + \mu(1+\rho)}.$$
 (A.2)

### **Real IR Function**

Here we present an expression for the IR function  $\tilde{B}$  which corresponds to the emission of a real photon  $k \in \Omega$  from a dipole consisting of two charged massive particles  $p_1$  and  $p_2$ .

$$2\alpha \tilde{B}(p_1, p_2) = -Z_i Z_j \theta_i \theta_j \frac{\alpha}{\pi} \bigg[ \bigg( \frac{1}{\rho} \ln \frac{\mu(1+\rho)}{m_1 m_2} - 1 \bigg) \ln \frac{\omega}{m_\gamma^2} + \frac{1}{2\beta_1} \ln \frac{1+\beta_1}{1-\beta_1} + \frac{1}{2\beta_2} \ln \frac{1+\beta_2}{1-\beta_2} + \mu G(p_1, p_2) \bigg],$$
(A.3)

logarithms and dilogarithms,

$$G(p_1, p_2) = \frac{1}{\sqrt{(Q^2 + M^2)(Q^2 + \delta^2)}} \left[ \ln \frac{\sqrt{\Delta^2 + Q^2} - \Delta}{\sqrt{\Delta^2 + Q^2} + \Delta} \left[ \chi_{23}^{14}(\eta_1) - \chi_{23}^{14}(\eta_0) \right] + Y(\eta_1) - Y(\eta_0) \right],$$
(A.4)

# **Progress on CEEX in Fortran**

[2203.10948] **Sherpa YFS** 

where  $\beta_i = \frac{|\vec{p}_i|}{E_i}$  and  $\mu, \rho$  are defined in Eq. (A.2) and  $\omega$  is the momentum cut-off specifying  $\Omega$  in the frame  $\tilde{B}$  is to be evaluated in.  $G(p_1, p_2)$  is a complicated function that can be expressed as a combination of

![](_page_20_Picture_22.jpeg)

See YFS.nb notebook in the NOTES folder of the git repository

![](_page_20_Figure_25.jpeg)

![](_page_21_Picture_0.jpeg)

![](_page_21_Picture_1.jpeg)

# ANALYSIS

# Modules HISTO, VARIABLES

Jérémy Paltrinieri

# **Progress on CEEX in Fortran**

# Progress on CEEX in Fortran

# ANALYSIS

# Modules HISTO, VARIABLES

This module allows to **create, fill, display, and save histograms at running time** (MC event generation).

### What each part does:

• init\_histo:

Sets up empty histogram bins between a lower and upper limit. Each bin gets a range and starts with a value of zero.

• fill\_histo:

Adds the *weight* of the event to the correct bin based on the input data. It figures out which bin the data point belongs to and increases that bin's count.

• finalize\_histo:

Finishes the histogram by adjusting all the values so they're properly scaled.

• print\_histo:

Simply prints out the histogram — showing the lower and upper edge of each bin and the value it holds.

 export\_histogram\_to\_csv: Saves the histogram to a CSV file generators.

## Jérémy Paltrinieri

Saves the histogram to a CSV file that can be compared to other MC

![](_page_22_Picture_17.jpeg)

# **Progress on CEEX in Fortran**

# ANALYSIS

# Modules HISTO, VARIABLES

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• export\_histogram\_to\_csv: generators.

## Jérémy Paltrinieri

Saves the histogram to a CSV file that can be compared to other MC

![](_page_23_Picture_17.jpeg)

Low         Up         Va           0.000000         0.200000         0.00000           0.200000         0.400000         0.00000           0.400000         0.600000         0.00000           0.600000         0.800000         0.00000           0.800000         1.000000         0.43349	Finalized Histogram				
0.0000000.2000000.000000.2000000.4000000.000000.4000000.6000000.000000.6000000.8000000.000000.8000001.0000000.000001.0000001.2000000.43349	Value				
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	D000E+00 D000E+00 D000E+00 D000E+00 D000E+00 D000E+00 D000E+00 D000E+00 D000E+00				

 $m_{\mu\mu}^{2}$  invariant mass of LO events

# Validation of Fixed-Order Modules

As is Phokhara, everything can be run in parallel, and of course for LO events the generation and integration is very fast.

## Jérémy Paltrinieri

![](_page_24_Figure_4.jpeg)

![](_page_24_Picture_5.jpeg)

# Validation of Fixed-Order Modules

As is Phokhara, everything can be run in parallel, and of course for LO events the generation and integration is very fast.

```
[jpaltrinieri@kappa CEEX]$ ./MC 1 10000000
RUNNING FORTRAN CEEX PROGRAM
Seed set to:
                1
         1000000
NbEvents:
PRINTING RESULTS
Monte Carlo Integral Result (pb):
                         86788.666725620918
Analytical Result Massive (pb):
                        86788.525413961281
Analytical Result Massless (pb):
                         86854.463232265436
Ratio = 1.0000016282297570
PROGRAM COMPLETE
```

## Jérémy Paltrinieri

![](_page_25_Figure_4.jpeg)

generator)

# Fortran modules in Phokhara

## $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ @ NNLO See William's talk

CEEX formalism implemented in Phokhara, with the goal of implementing resummed  $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ 

valid at  $\mathcal{O}(\alpha^1)_{CEEX}$  in the language of KKMC

Jérémy Paltrinieri

![](_page_26_Picture_6.jpeg)

# Fortran modules in Phokhara

## $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ @ NNLO See William's talk

CEEX formalism implemented in Phokhara, with the goal of implementing resummed  $e^+e^- \rightarrow \mu^+\mu^-\gamma, \pi^+\pi^-\gamma$ 

valid at  $\mathcal{O}(\alpha^1)_{CEEX}$  in the language of KKMC

Jérémy Paltrinieri

All modules implemented for the CEEX resummation framework can be added to the Phokhara framework itself. There will be a switch between the fixed-order and the resummed calculation

![](_page_27_Picture_7.jpeg)

# **Progress on CEEX in Phokhara**

![](_page_28_Figure_1.jpeg)

## Summary

- Framework in Fortran established
- Fixed-Order Modules working
- Debugging multi-photon Modules

Jérémy Paltrinieri

![](_page_28_Picture_9.jpeg)

# **Progress on CEEX in Phokhara**

![](_page_29_Picture_1.jpeg)

# Summary

- Framework in Fortran established
- Fixed-Order Modules working
- Debugging multi-photon Modules -----

![](_page_29_Picture_6.jpeg)

# What's next?

- Validation of KKMC-like exponentiation
- Incorporation in Phokhara -
- -----
- GVMD@NLO for Pions, but no plans for GVMD@NNLO

## Jérémy Paltrinieri

Improving the beta functions, rewrite the procedure in dim-reg